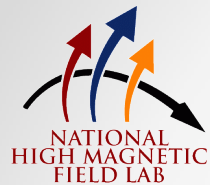


Shallow Donors in Silicon as Electron and Nuclear Spin Qubits

Johan van Tol
National High Magnetic Field Lab
Florida State University

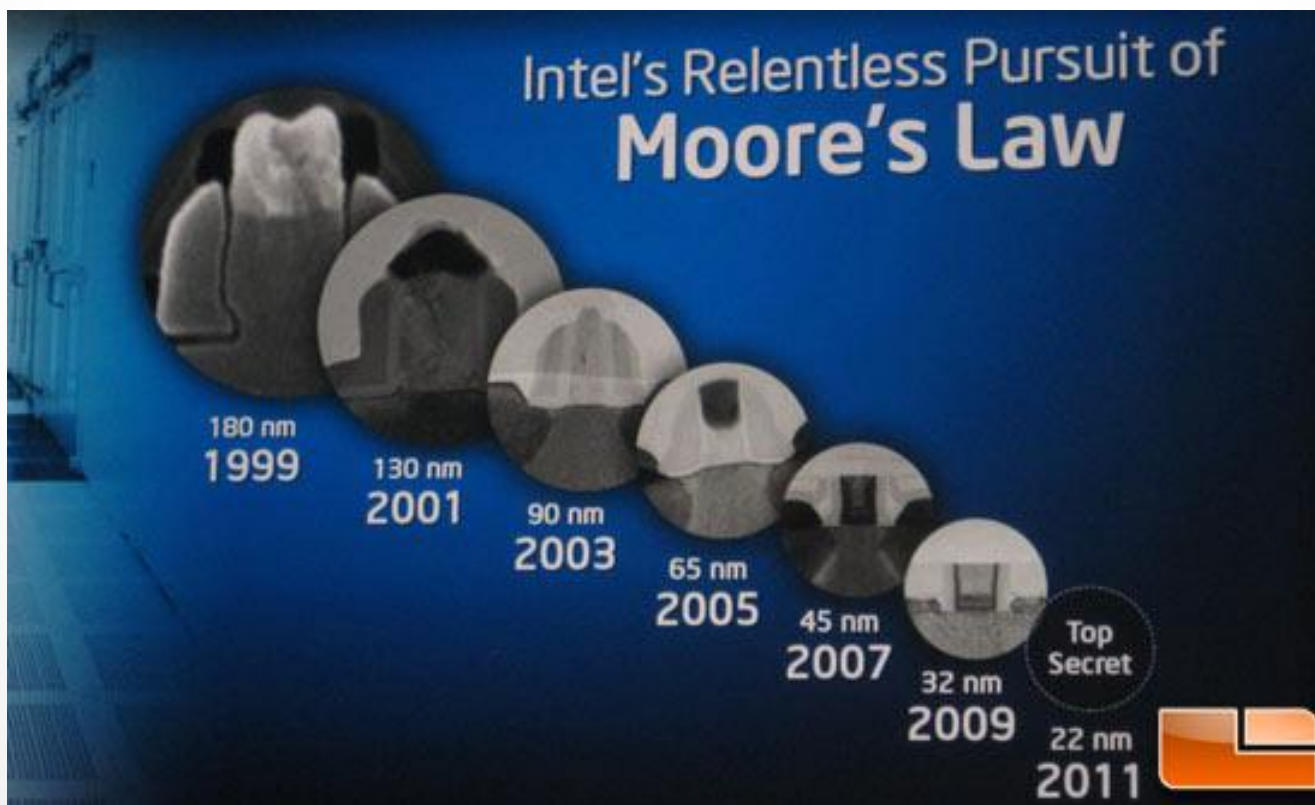


Overview

- Electronics – The end of Moore's law?
- Quantum computing
- Spin qubits
- Relaxation
- Manipulation
- Readout
- FEL EPR

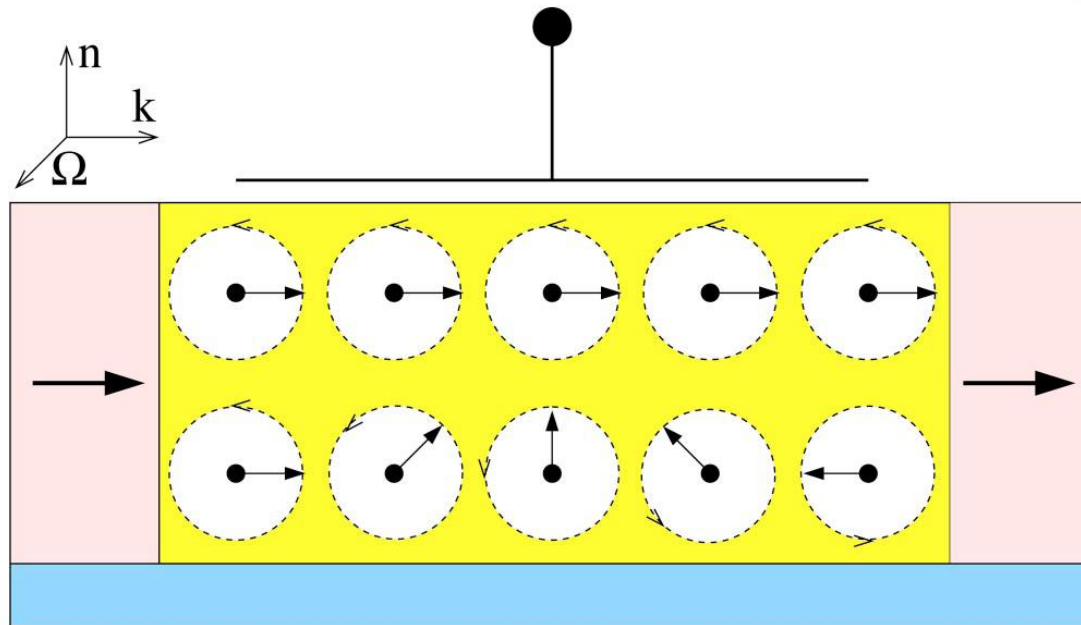
(Gordon) Moore's law

- "The number of transistors incorporated in a chip will approximately double every 24 months."



Spintronics

- Electronics depends on the charge of the electron
- Using the spin: **spin** transport electronics



Source: Zutic et al. *Rev Mod Phys* 76 (2004), 323.

Datta-Das Spin Field Effect Transistor

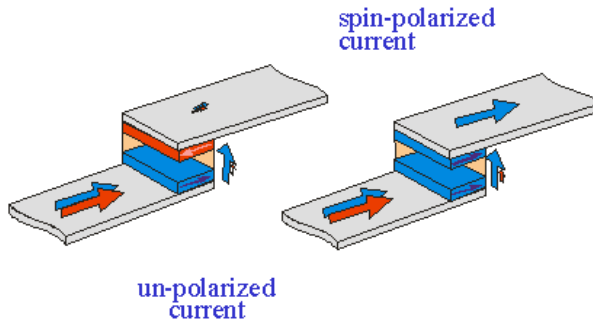
Spintronics device MRAM

- Universal Memory

MTJ

*insulating space layer
current perpendicular to
plane*

Ferromagnetic
electrode 1
Tunneling
barrier
Ferromagnetic
electrode 2

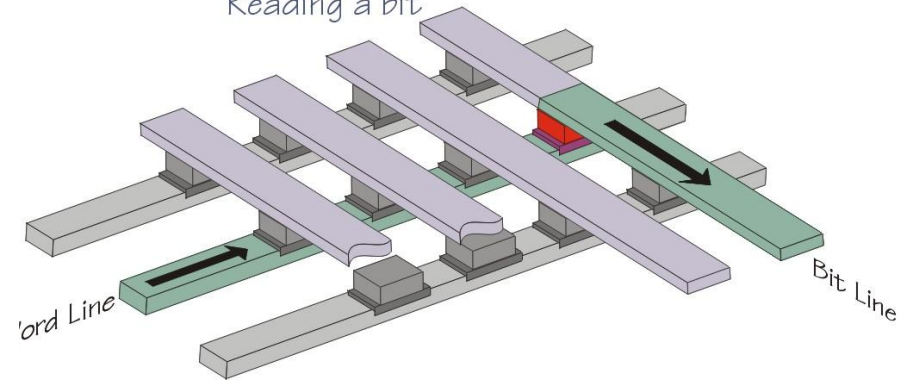


first ferromagnetic electrode acts as spin filter
second FM layer acts as spin detector

S.S.P. Parkin
December 16, 1998

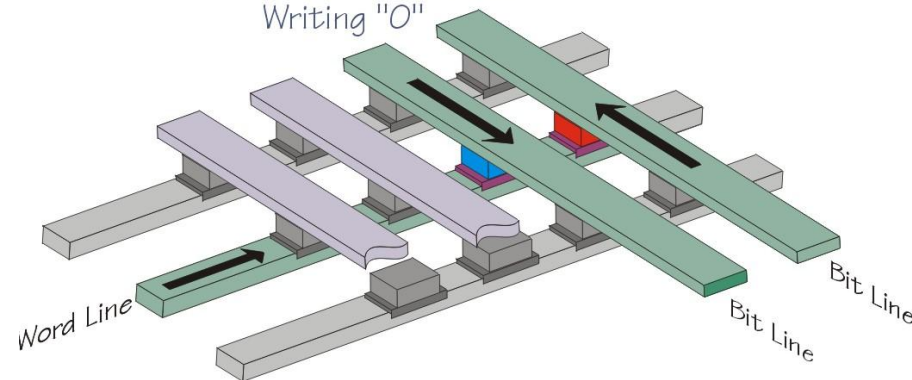
MagRAM Architecture

Reading a bit



Writing "1"

Writing "0"



MTJ MagRAM promises

- density of DRAM
- speed of SRAM
- non-volatility



Quantum Computer

Simulating Physics with Computers

Richard P. Feynman

Department of Physics, California Institute of Technology, Pasadena, California 91107

Received May 7, 1981

1. INTRODUCTION

On the program it says this is a keynote speech—and I don't know what a keynote speech is. I do not intend in any way to suggest what should be in this meeting as a keynote of the subjects or anything like that. I have my own things to say and to talk about and there's no implication that anybody needs to talk about the same thing or anything like it. So what I want to talk about is what Mike Dertouzos suggested that nobody would talk about. I want to talk about the problem of simulating physics with computers and I mean that in a specific way which I am going to explain.

Peter Shor:

Factorization of a prime number can be solved in many fewer steps in a quantum computer. (1994)



Requirements

- Scalable quantum system that:
 - Can represent quantum information robustly (long-lived: long relaxation times)
 - Perform transformations (operations)
 - Prepare an initial state
 - Measure the results

→ SPINS

Magnetic moment of electron and nuclear spins
is 'isolated' from the surroundings

Proposed schemes

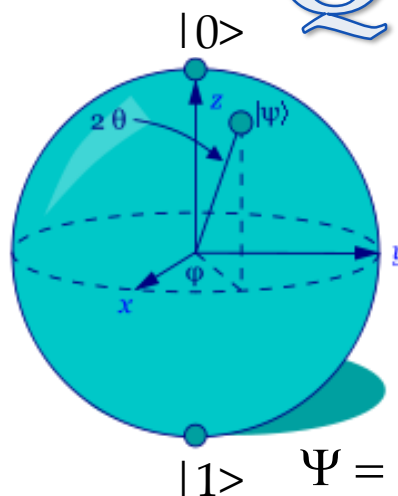
- [Superconductor](#)-based quantum computers (including [SQUID](#)-based quantum computers)^[18]
- [Trapped ion quantum computer](#)
- [Optical lattices](#)
- [Topological quantum computer](#)^[19]
- [Quantum dot](#) on surface (e.g. the [Loss-DiVincenzo quantum computer](#))
- [Nuclear magnetic resonance](#) on molecules in solution (liquid NMR)
- Solid state NMR [Kane quantum computers](#)
- Electrons on helium quantum computers
- [Cavity quantum electrodynamics](#) (CQED)
- [Molecular magnet](#)
- [Fullerene](#)-based [ESR](#) quantum computer
- Optic-based quantum computers ([Quantum optics](#))
- [Diamond-based quantum computer](#)^{[20][21][22]}
- [Bose–Einstein condensate-based quantum computer](#)^[23]
- Transistor-based quantum computer - string quantum computers with entrainment of positive holes using an electrostatic trap
- [Spin-based quantum computer](#)
- [Adiabatic quantum computation](#)^[24]
- Rare-earth-metal-ion-doped inorganic crystal based quantum computers^{[25][26]}

Quantum Computing

Bit

1 or 0

QuBit



$$|\Psi\rangle = \cos(\theta)|0\rangle + \exp^{-i\phi}\sin(\theta)|1\rangle$$

2 bits: 4 states: 00, 01, 10, 11

2 qubits: $|00\rangle$, $|01\rangle$, $|10\rangle$, $|11\rangle$,

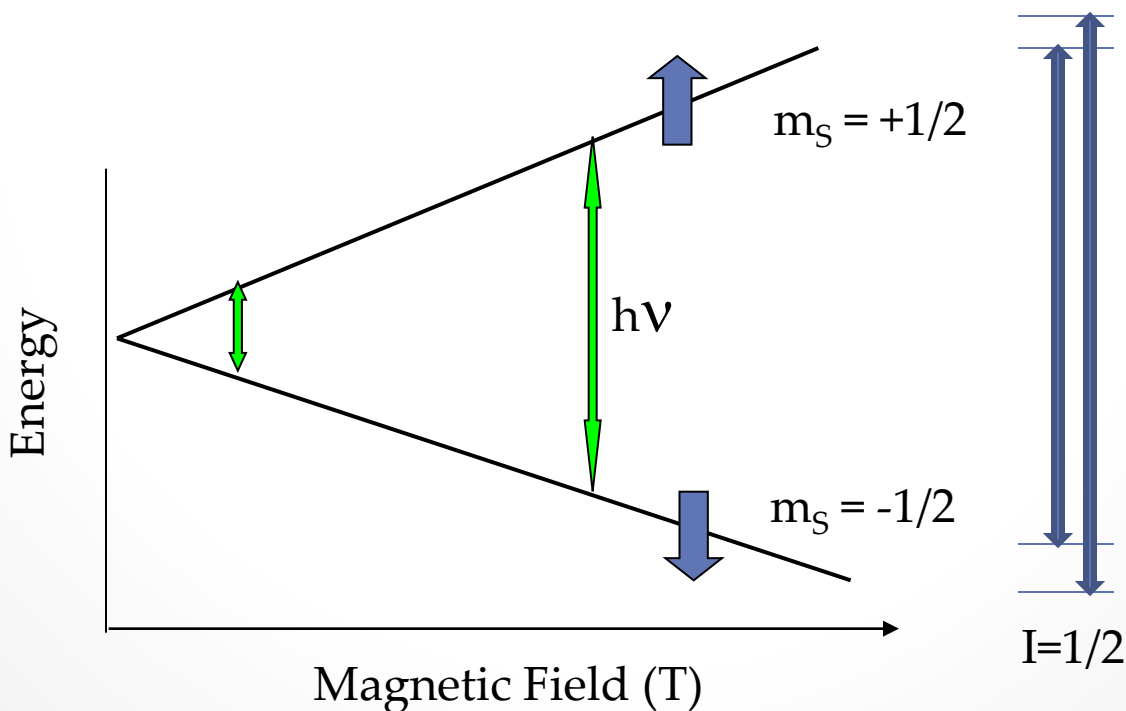
$\frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$, $\frac{1}{\sqrt{2}} (|00\rangle + |01\rangle)$, ...

$\frac{1}{2} (|00\rangle - |01\rangle + |10\rangle - |11\rangle)$, ...



The EPR phenomenon

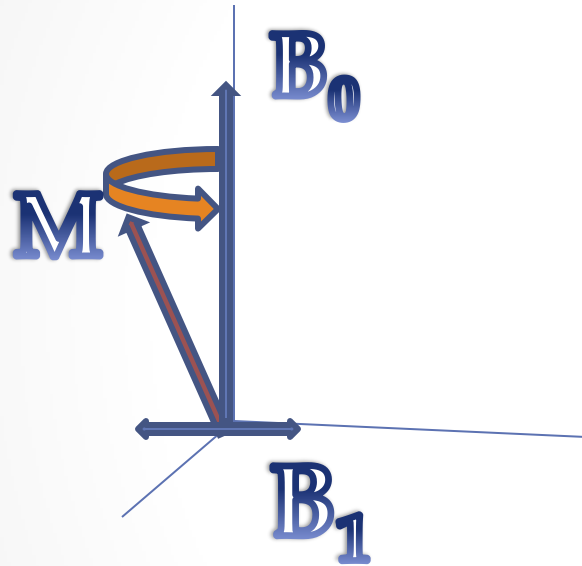
$$H = \underbrace{\mu_B B \cdot \mathbf{g} \cdot \mathbf{S}}_{\text{Zeeman}} + \underbrace{\mathbf{S} \cdot \mathbf{D} \cdot \mathbf{S}}_{\text{Fine Structure}} + \underbrace{\sum_j \mathbf{S} \cdot \mathbf{A} \cdot \mathbf{I}_j}_{\text{Hyperfine structure}} - \underbrace{\gamma_n I_j B}_{\text{Nuclear Zeeman}} + \underbrace{\sum_k J \mathbf{S}_k \cdot \mathbf{S}}_{\text{Exchange}} + \underbrace{\mathbf{S}_k \cdot \mathbf{T}_{\text{dip}} \cdot \mathbf{S}}_{\text{Dipolar}}$$



Change the reference frame



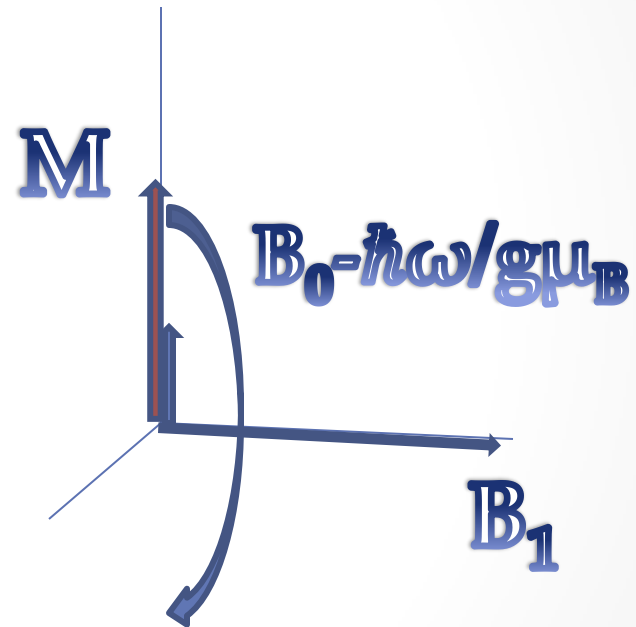
- In the lab frame



$$\frac{d\mathbf{M}(t)}{dt} = \gamma \mathbf{M}(t) \times \mathbf{B}(t)$$

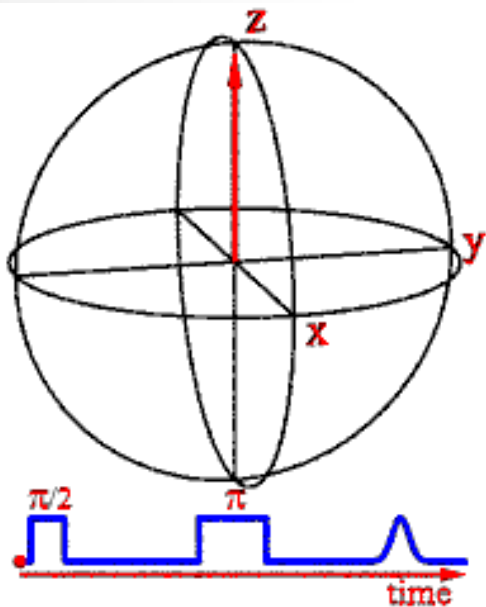
Precession around B_0

- In the rotating frame (ω)



Precession around B_1

What is pulse EPR/NMR ?



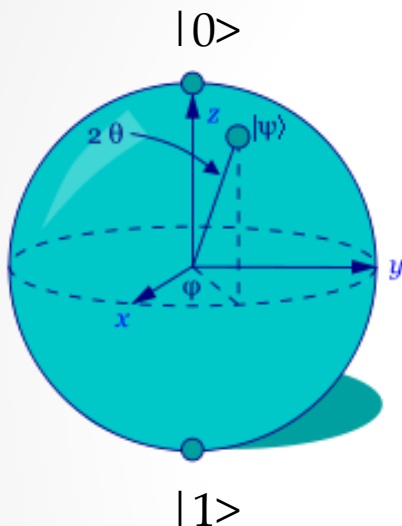
A spin echo seen in the rotating frame

Relaxation parameters

T_1 Spin-lattice relaxation time

T_2 Spin-memory time, spin-spin relaxation time

Relaxation



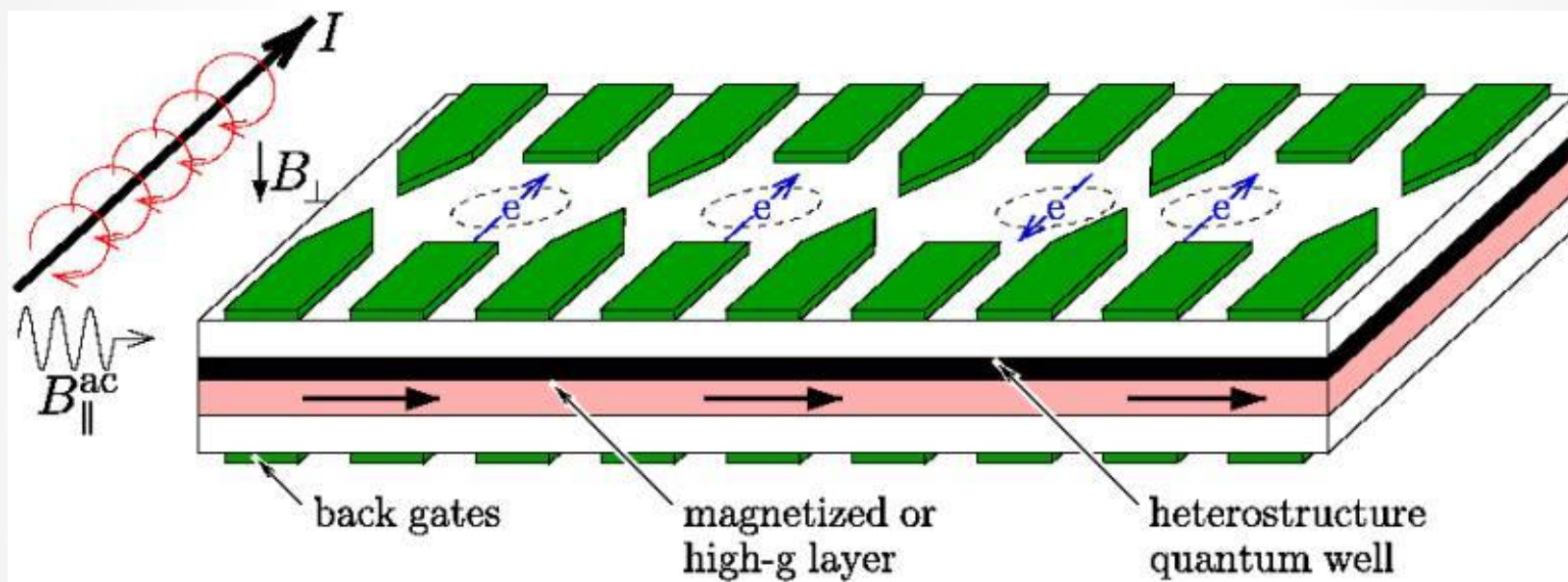
$$\Psi = \cos(\theta)|0\rangle + \exp^{-i\phi}\sin(\theta)|1\rangle$$

Loss of θ : T_1 (Energy)
Spin Lattice Relaxation
Longitudinal Relaxation

Loss of ϕ : T_2 (Phase)
Spin-Spin relaxation
Phase relaxation
Spin memory time
Coherence time

Quantum Dot QC

Loss and DiVincenzo, PRB '98

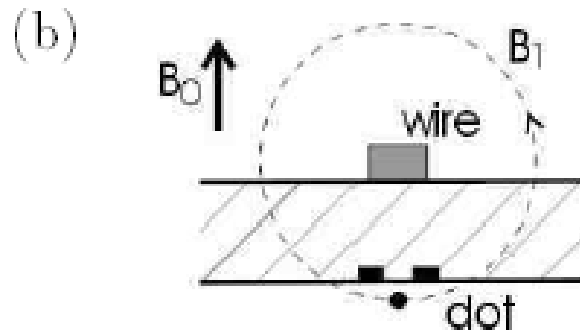
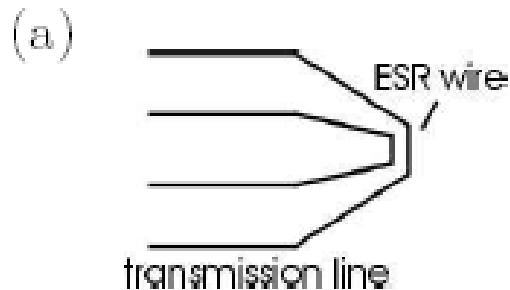
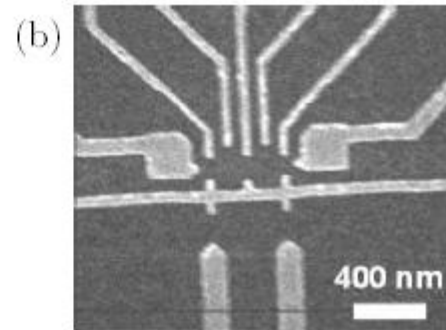
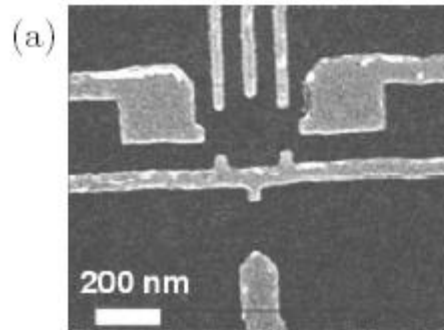


Elzerman, Kouwenhoven et al. (Delft University of Technology)

“Single-shot read-out of a Spin Qubit”

Electron Spin Resonance in quantum dots

Vandersypen et al. "Quantum Computing and Quantum Bits in Mesoscopic Systems", Kluwer



Shallow donors



Periodic Table of Elements

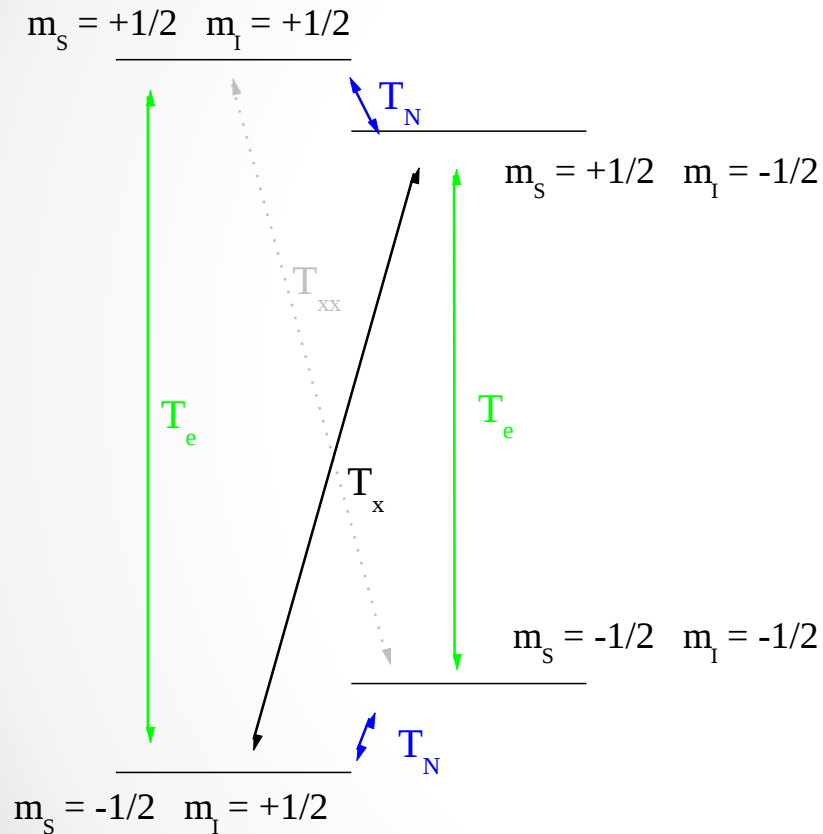
Periodic Table of Elements

										IIIA	IVA	VA	VIA	VIIA			
2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr			
Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe			
Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn			
Rf	Ha	106	107	108	109	110											

Si:P



- Spin system



$$H = gu_B \vec{S} \cdot \vec{B} - \gamma_n \vec{B} \cdot \vec{I} + a \vec{S} \cdot \vec{I}$$

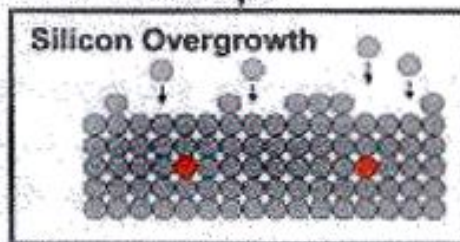
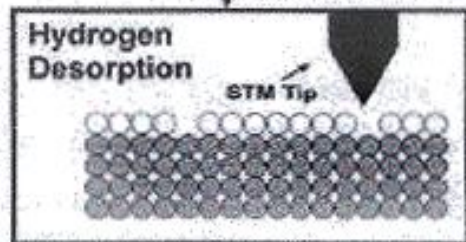
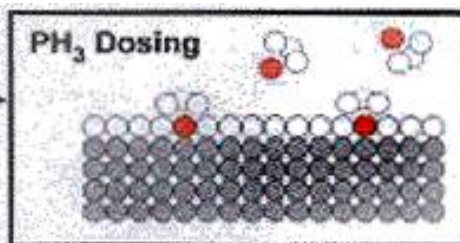
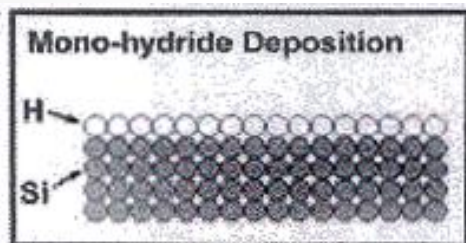
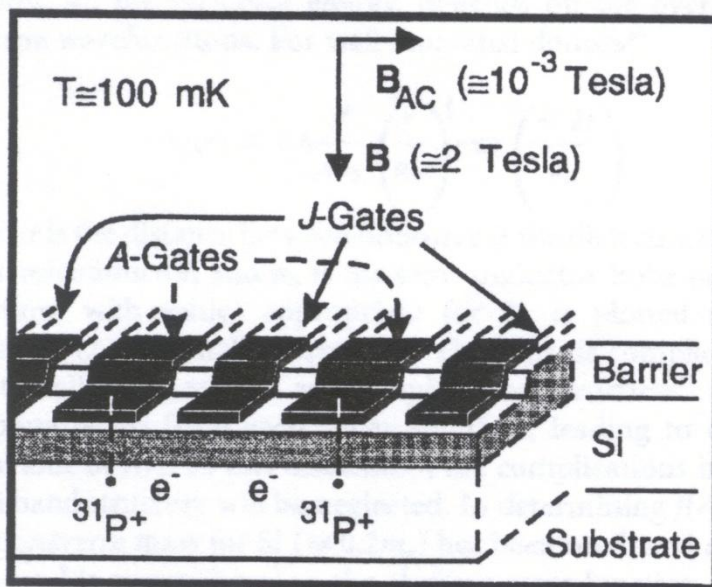
$$S_z I_z + \frac{1}{2} (S^+ I^- + S^- I^+)$$

+ $^{29}\text{Si} \dots$

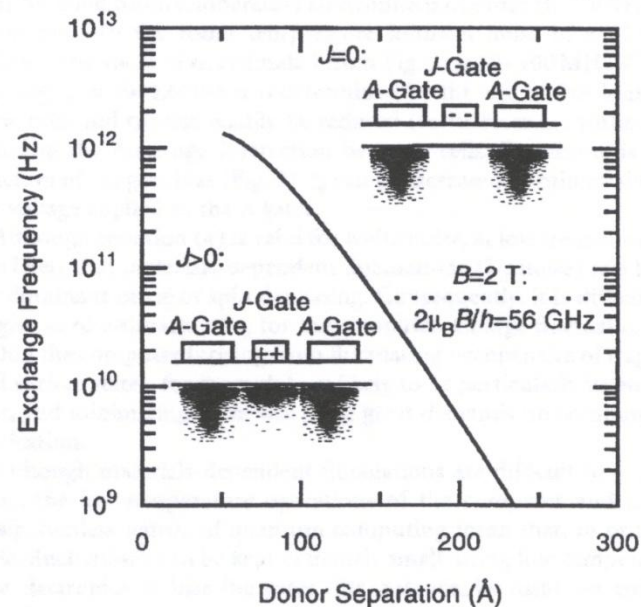
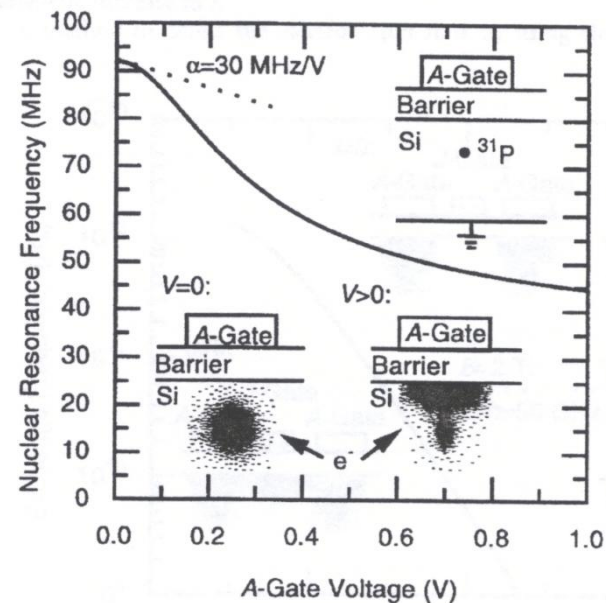
4 CNOT gates

Universal set of quantum gates

Kane's model

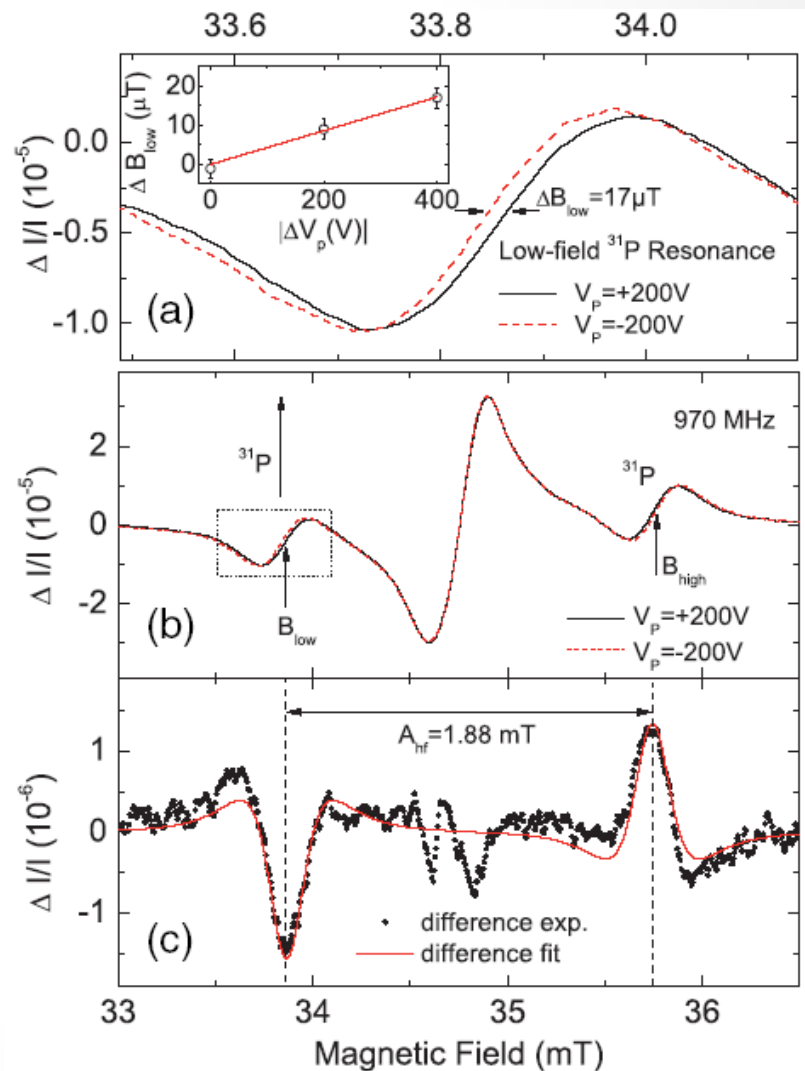
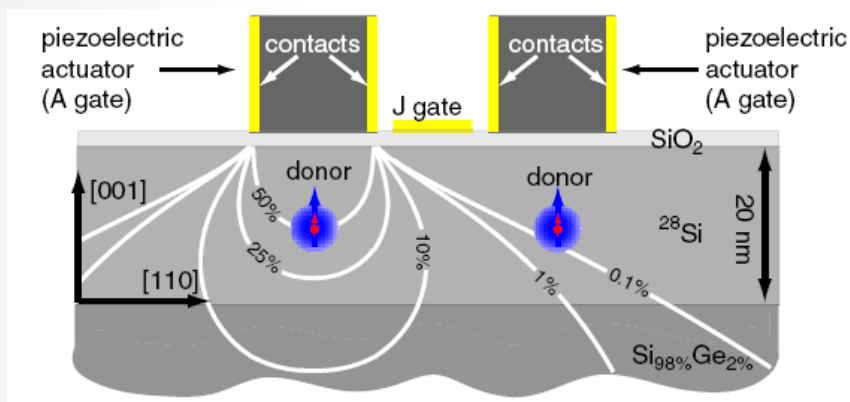


Nature, **393**, 133 (1998).

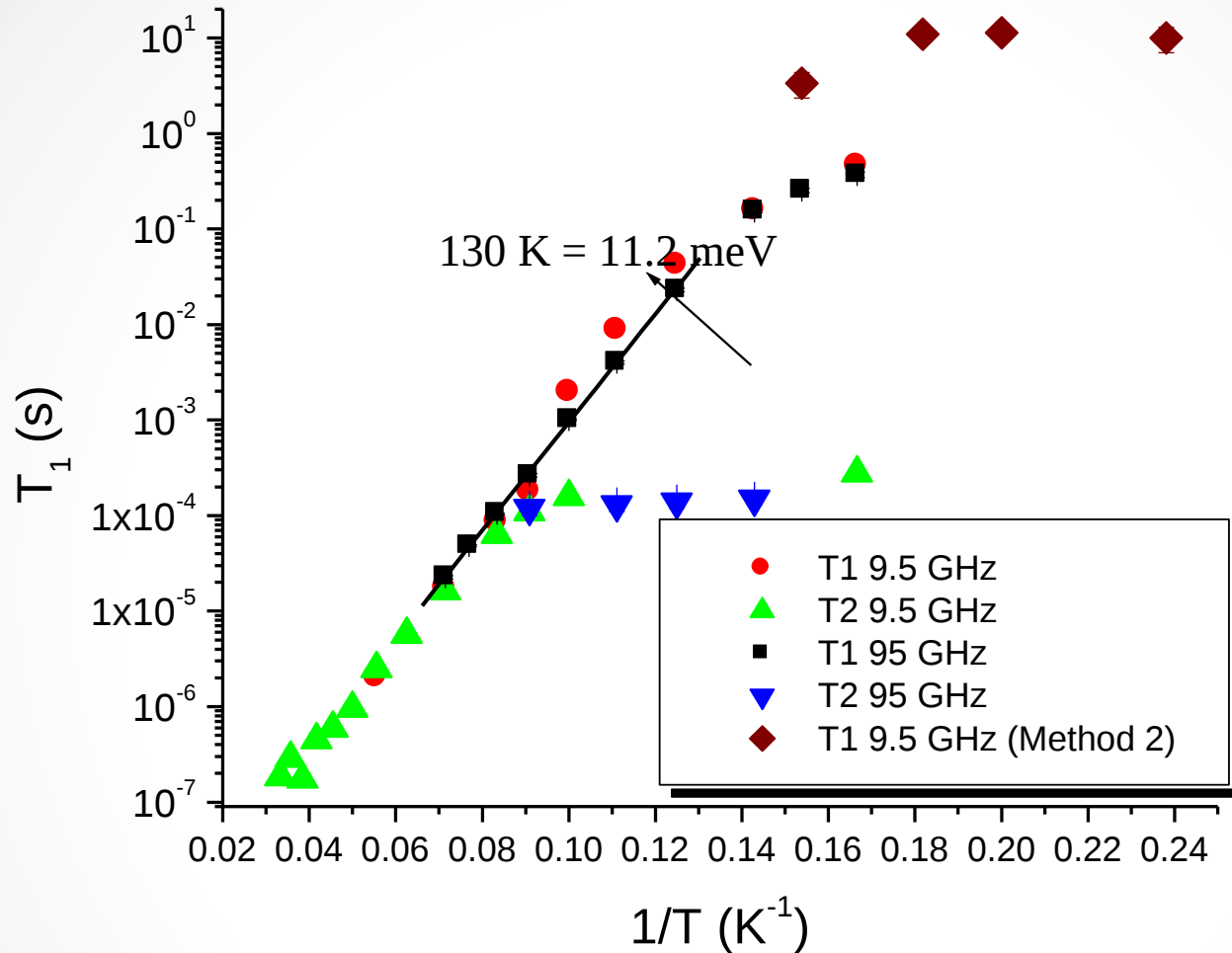


Tuning the hyperfine

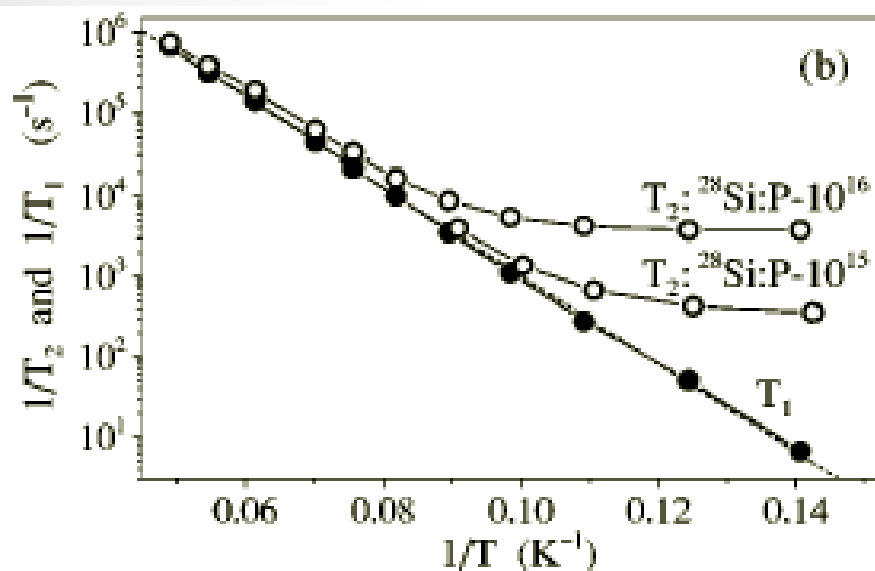
- Dreher et al. PRL 106, 037601 (2011)



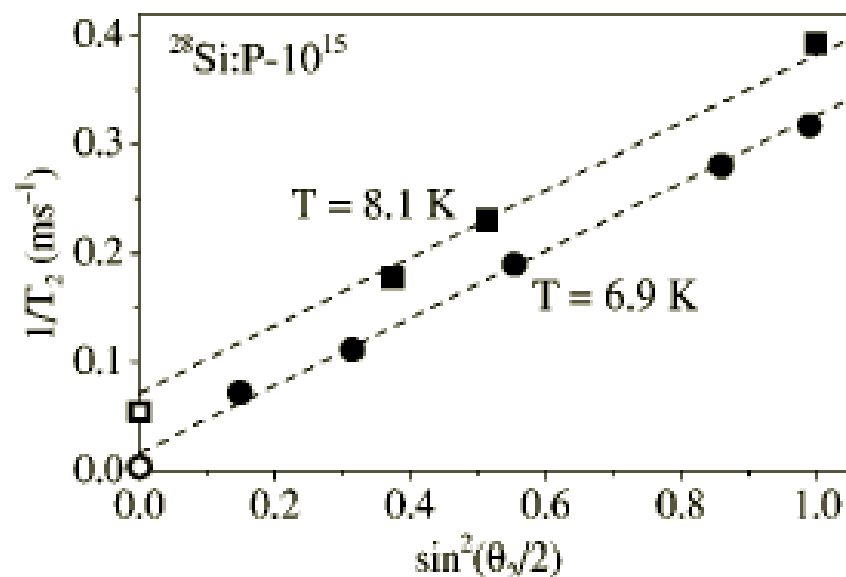
Electron Spin relaxation



- T_1 exponential temp. dependence: shallow donor excited state
- T_2 is constant ~ 200 - $300 \mu s$ due to ^{29}Si
- No field/frequency dependence up to 95 GHz down to 7 K

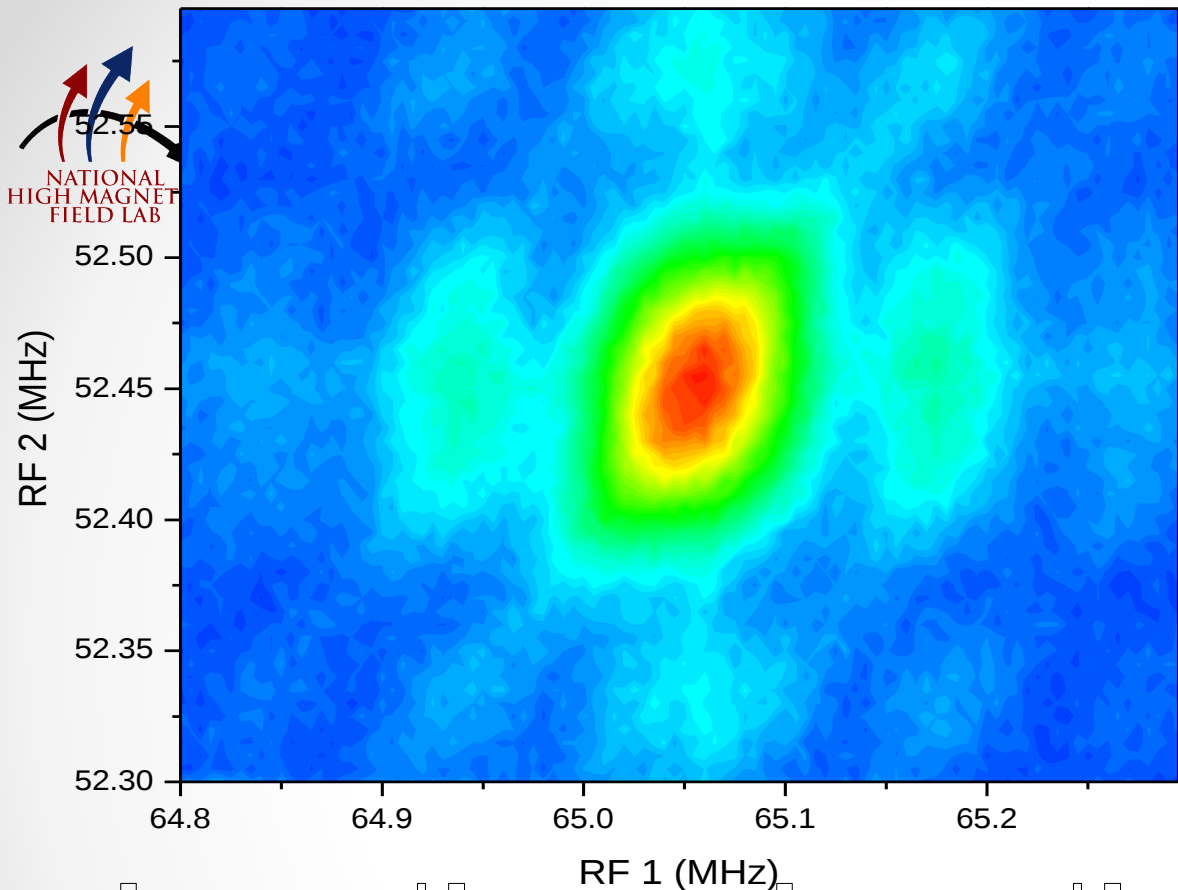


Temperature dependence T_1 , T_2



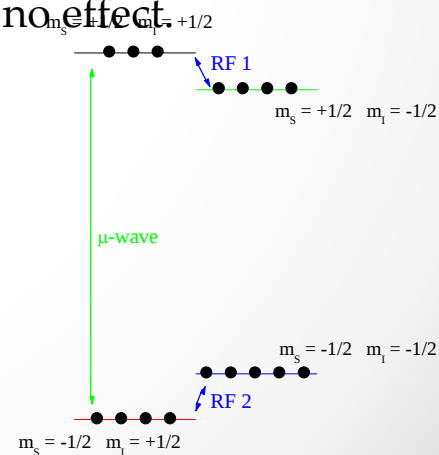
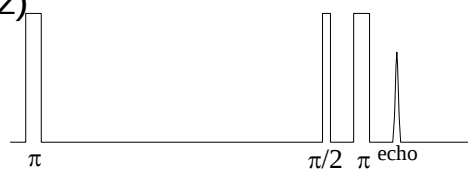
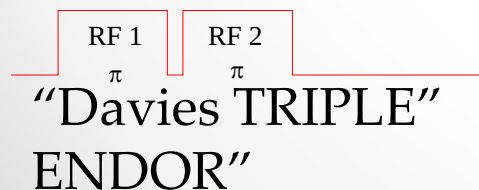
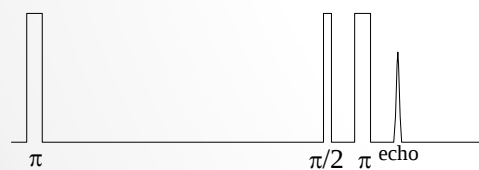
Instantaneous diffusion contribution

- In ^{28}Si samples T_{2e} is limited by dipolar interactions



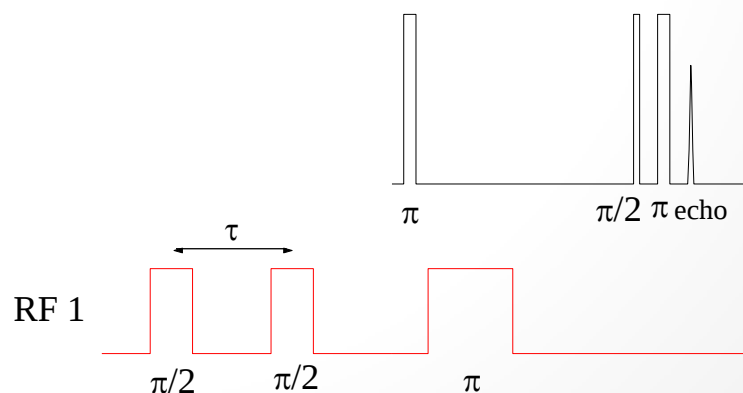
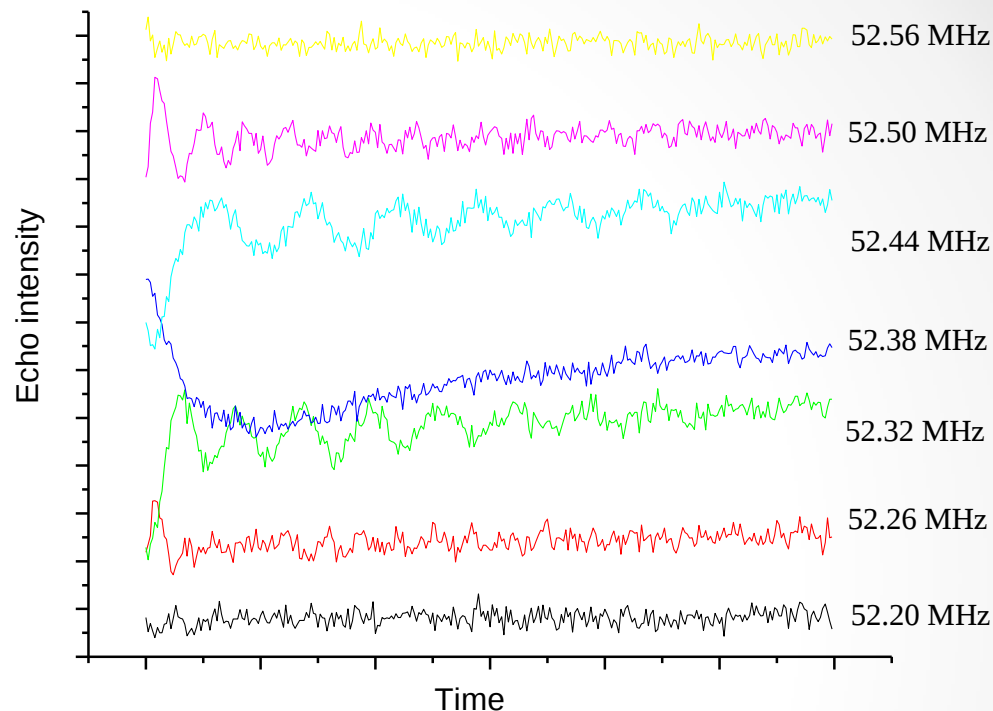
Strong TRIPLE enhancement

The pulsed ^{31}P TRIPLE spectrum in both Davies and Mims ENDOR is much stronger than the pulsed ENDOR itself, due to the slow relaxation of the nuclear spins. After a typical Davies-ENDOR sequence, followed by electron-spin T_1 decay, the populations of the levels end up as shown below – after the electron spin inversion pulse, the RF pulse has no effect.

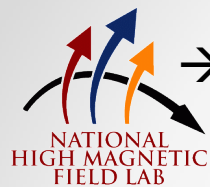


^{31}P free induction decay, measured by “ENDOR”

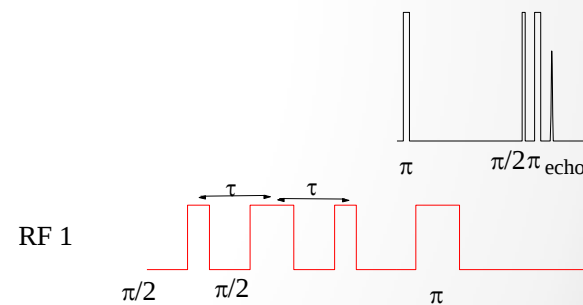
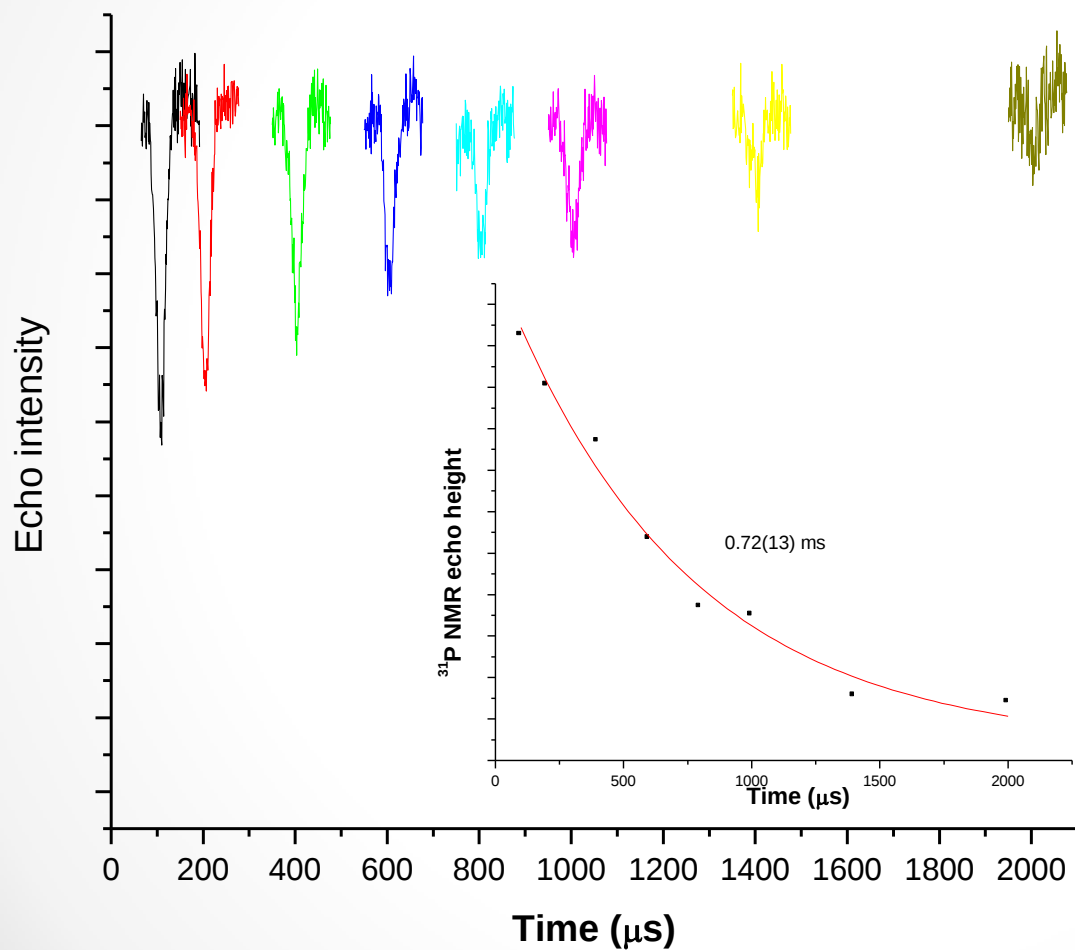
The pulsed-ENDOR sequence can be used to both induce the nuclear polarization and to detect the NMR signal of the ^{31}P nuclei. The repetition rate is slow with respect to the electron spin T_1 , but fast with respect to the nuclear T_1 . The first RF pulse induces the free-induction decay, the second translates the nuclear coherence to a population difference, and is detected with the pulsed ENDOR sequence



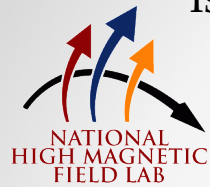
"NMR" echo, detected on EPR/ENDOR signal (10 K)



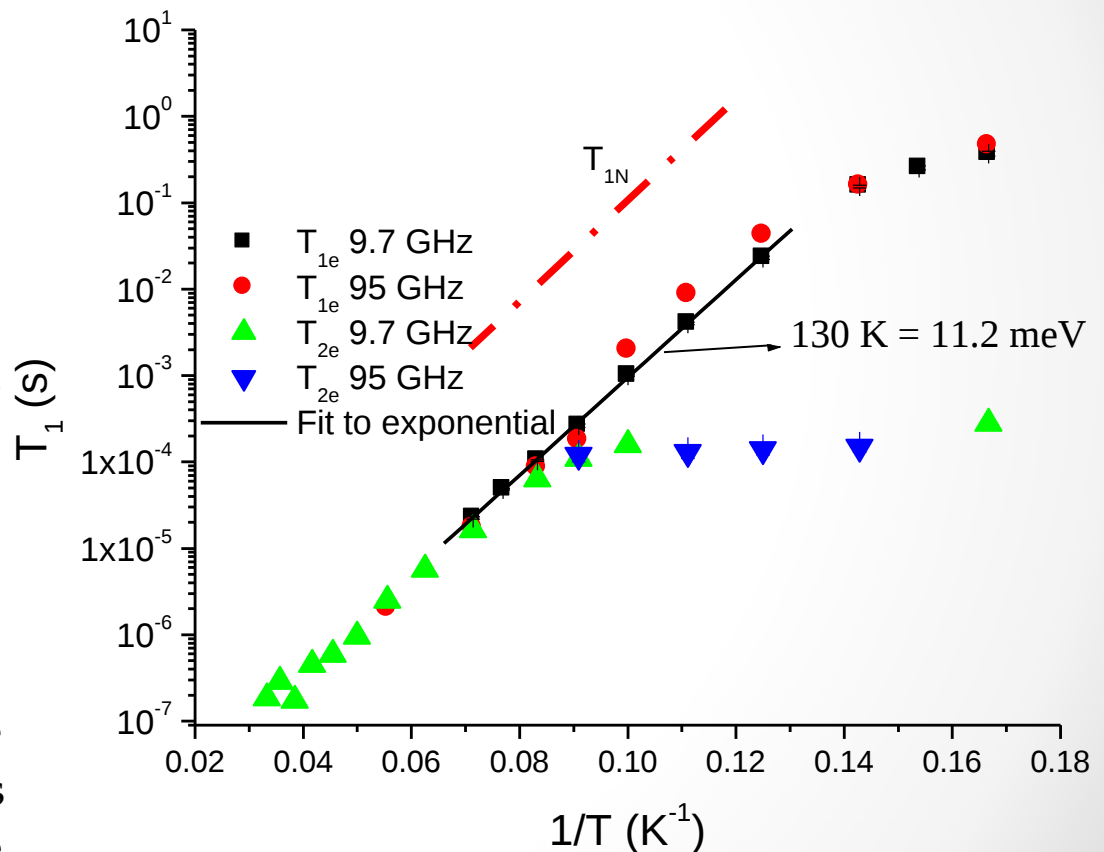
→ Nuclear coherence time T_{2N} seems to be close to T_{1e}



Si:P ([P] $\sim 10^{15}$ - 10^{16}) Below 20 K no conduction electrons: Nicely isolated from surroundings
Long relaxation times



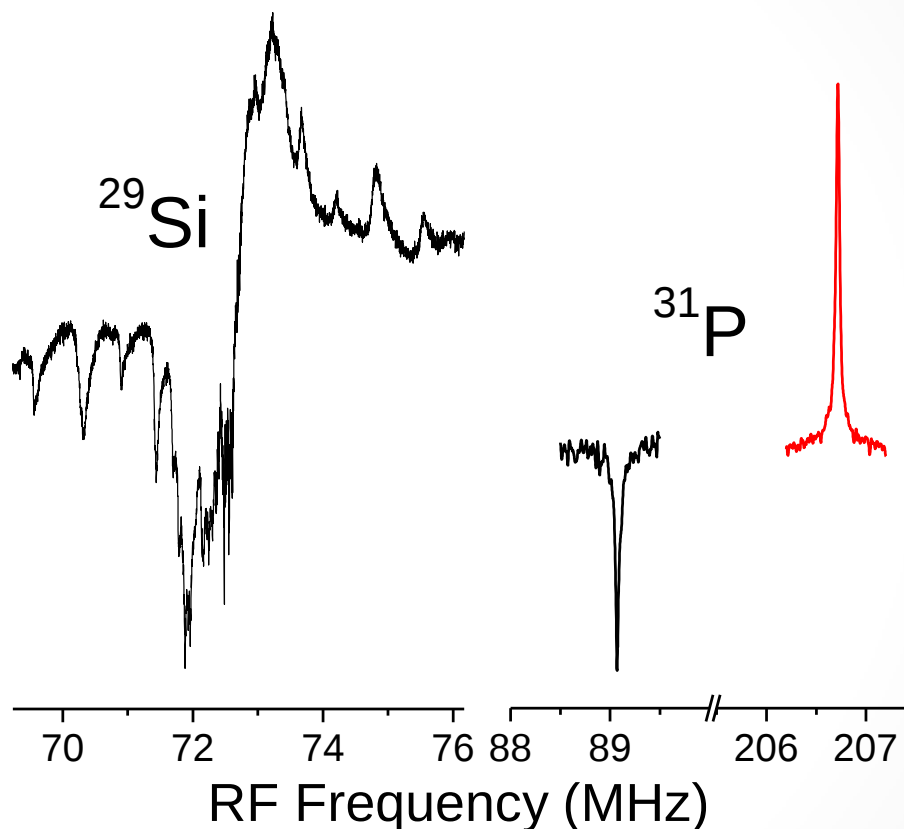
- T_{1e} exponential temp. dependence: shallow donor excited state
- T_{2e} is constant ~ 200 - $300 \mu\text{s}$ due to ^{29}Si
- No field/frequency dependence up to 95 GHz
- The ^{31}P nuclear T_{1N} relaxation rate at these temperatures and at 0.35 T is about 2 orders of magnitude slower than T_{1e} , while the nuclear T_{2N} of the phosphorus spins is limited by the spin-lattice relaxation rate of the electron spins ($T_{2N} \sim T_{1e}$)



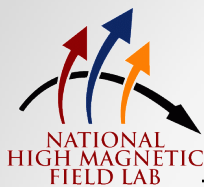
240 GHz ENDOR

^{29}Si ENDOR from different shells. Strongly coupled Si relax much slower than the weakly coupled Si.

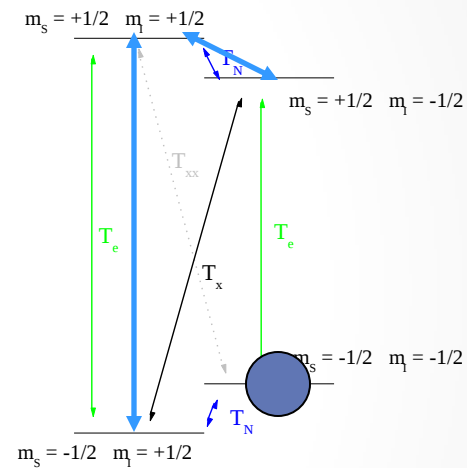
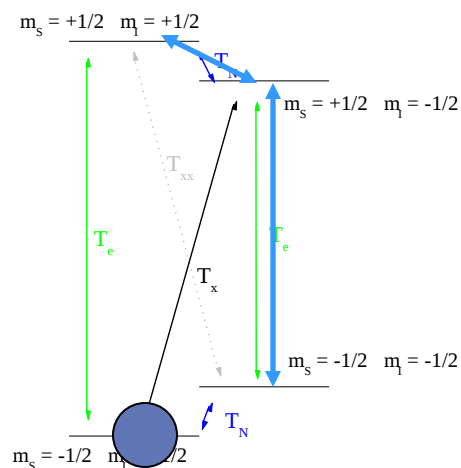
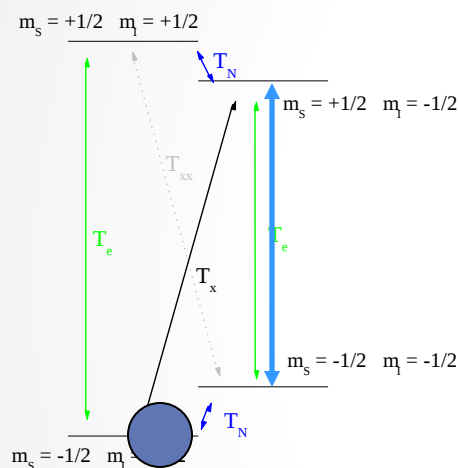
In principle they could add another qubit to the system.



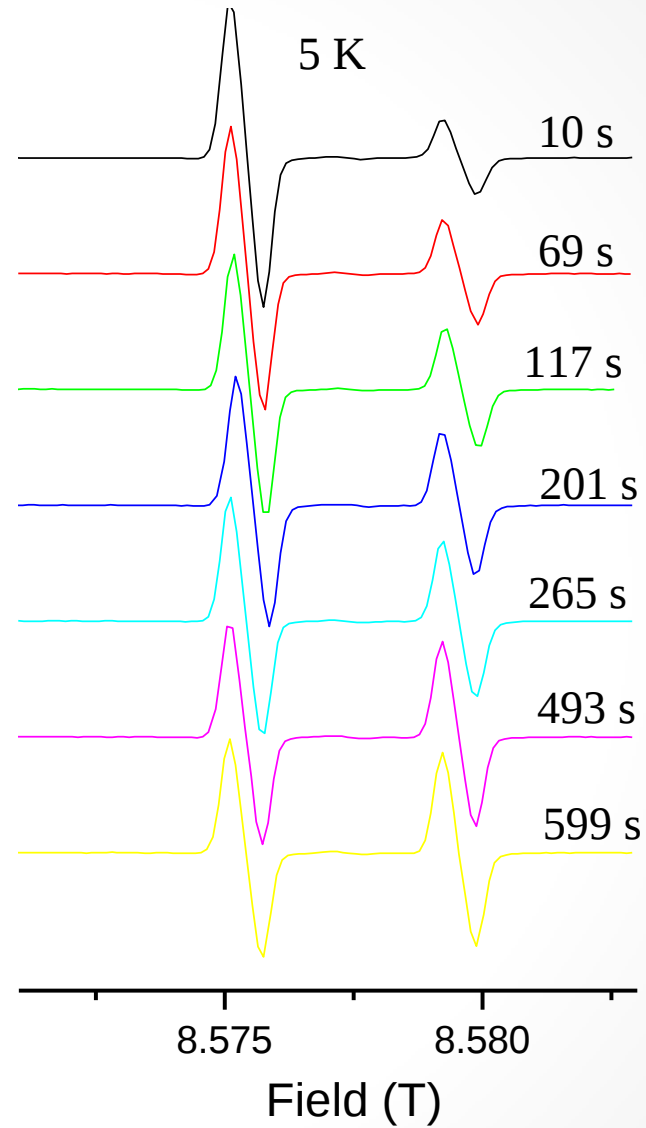
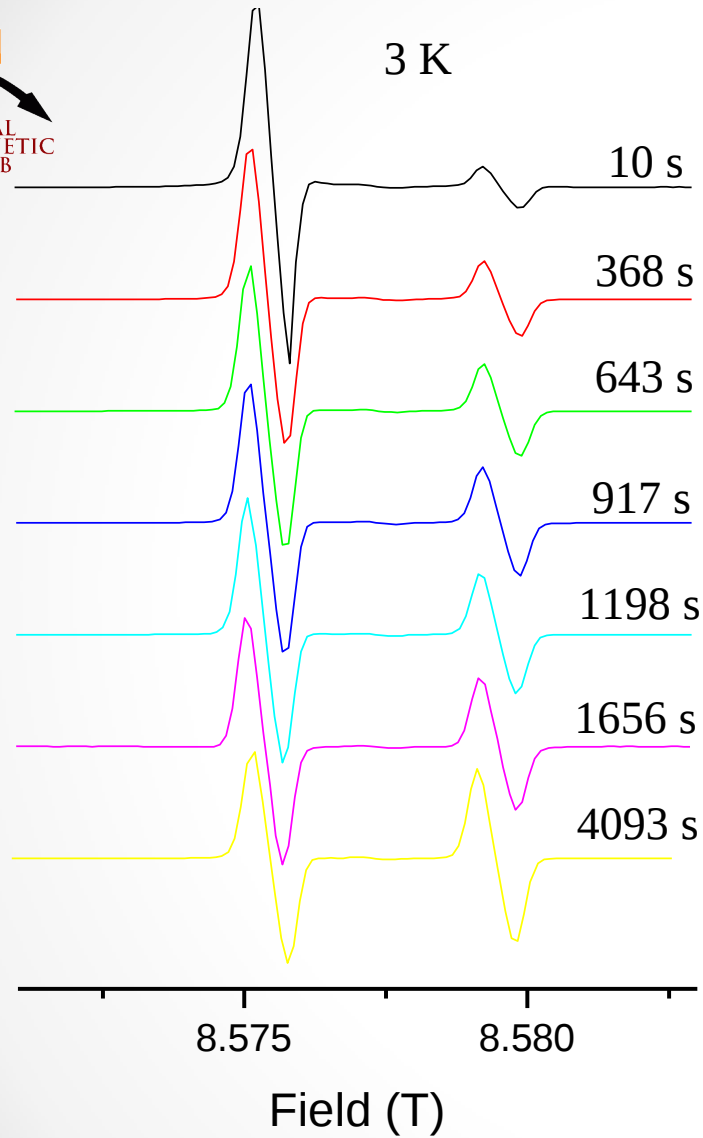
High-field dynamic nuclear polarization at $h\nu \gg kT$

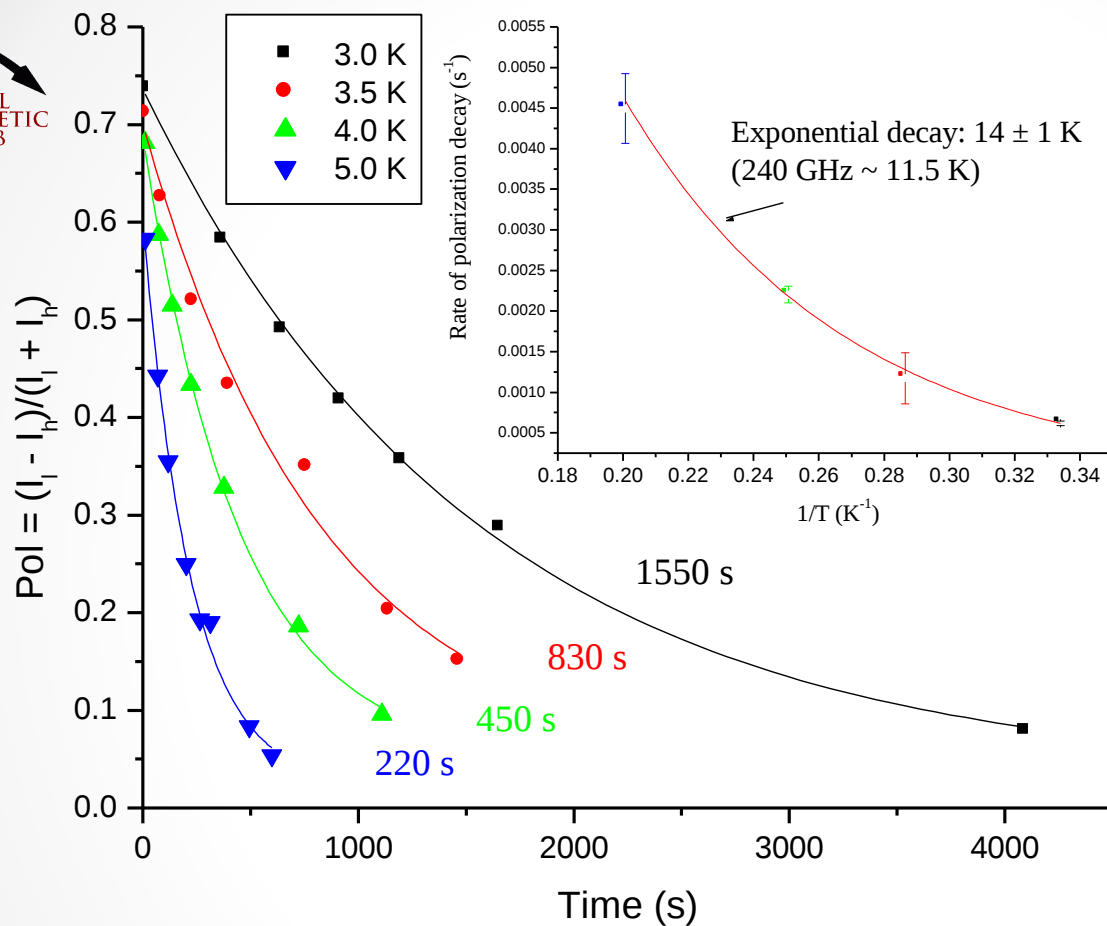


Creation of close to 100% polarized initial state



Initialization/Relaxation of nuclear spin polarization





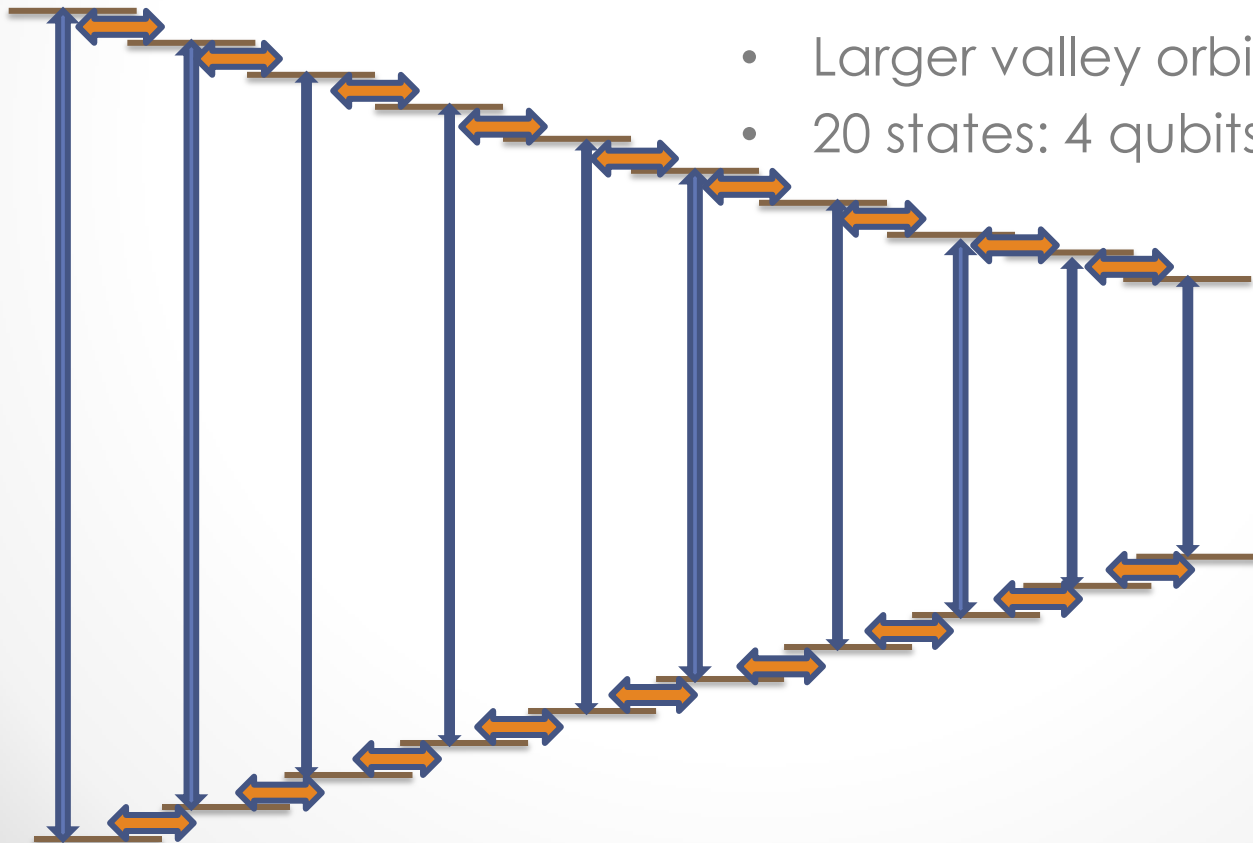
The relaxation to equilibrium spin polarization is well described with a single exponential. The temperature dependence indicates a thermally excited process with an energy close to the electron spin Zeeman splitting at these fields.

No polarization is achieved by saturation of the low-field hyperfine component.

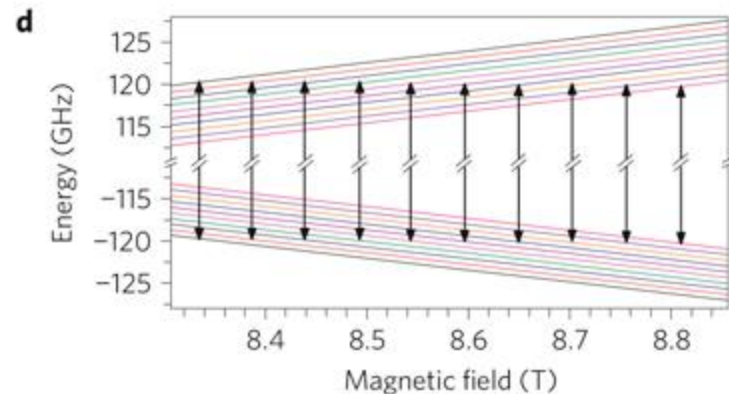
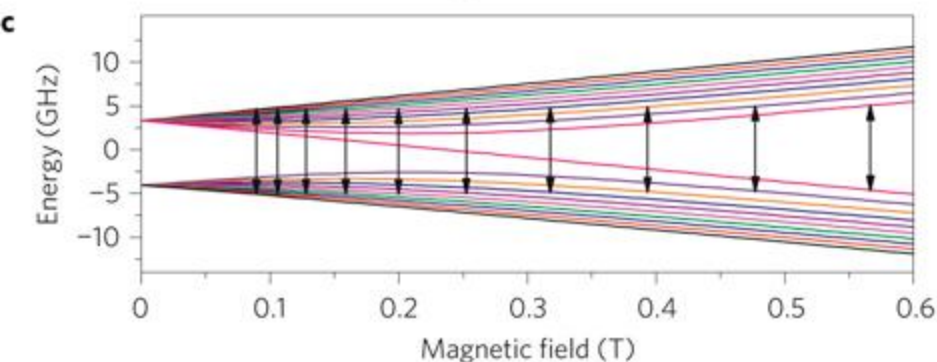
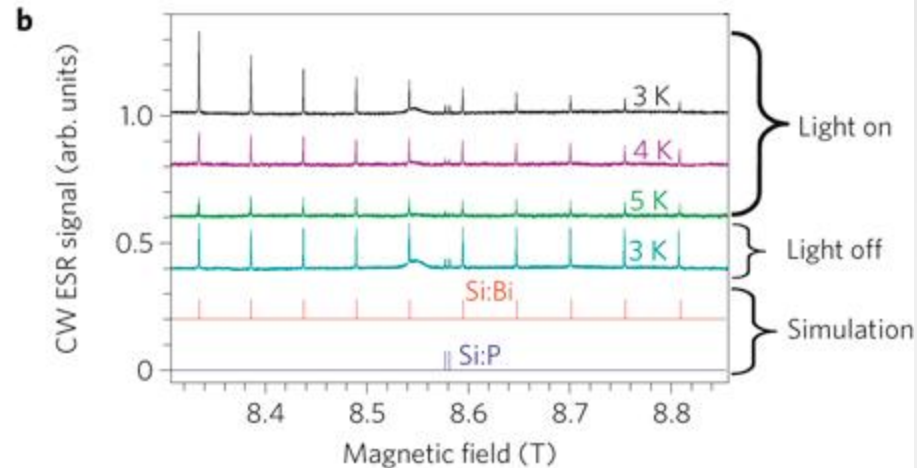
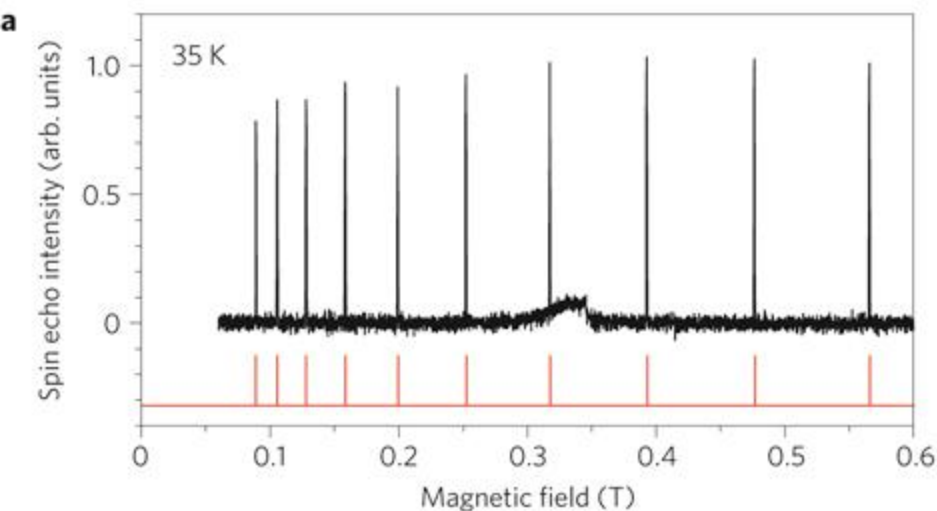
Bismuth in Silicon



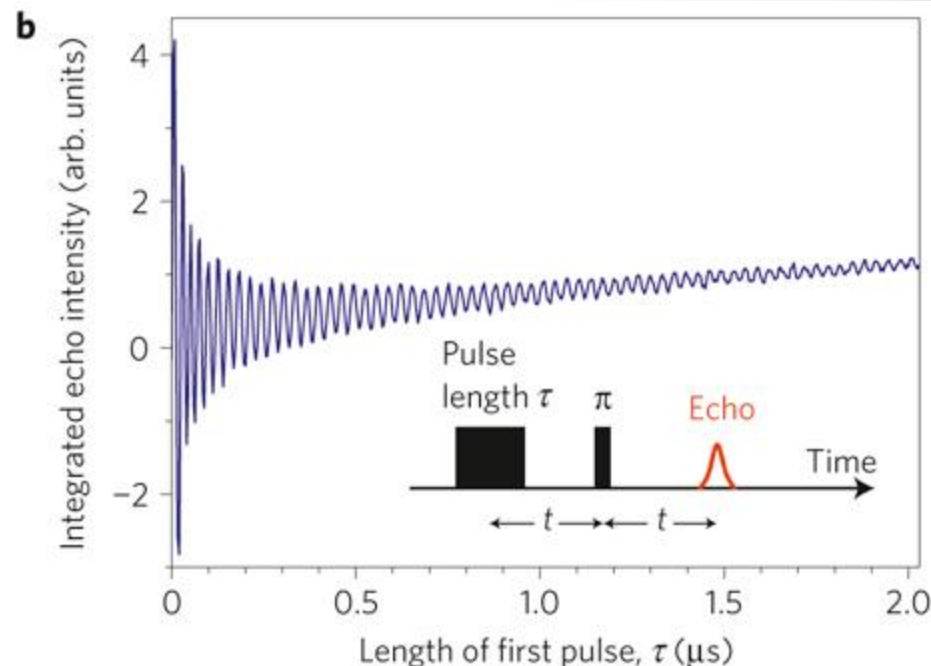
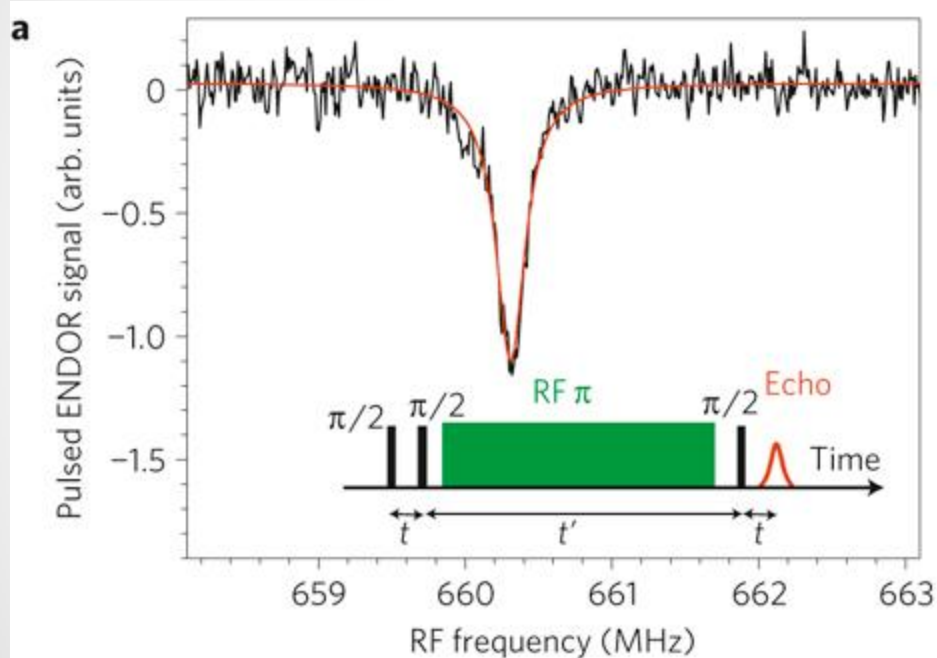
- $I=9/2$
- Relaxation times similar to phosphorus
- Much larger hyperfine splitting $\rightarrow$ individual addressing
- Larger valley orbit splitting
- 20 states: 4 qubits at once ?



Another candidate: Bi in Si



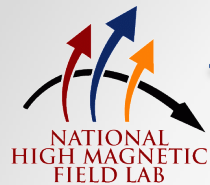
Electron Nuclear Double Resonance and coherent manipulation



Bismuth is the heaviest stable element.

It has a large nuclear spin $9/2 \rightarrow$ More information can be stored

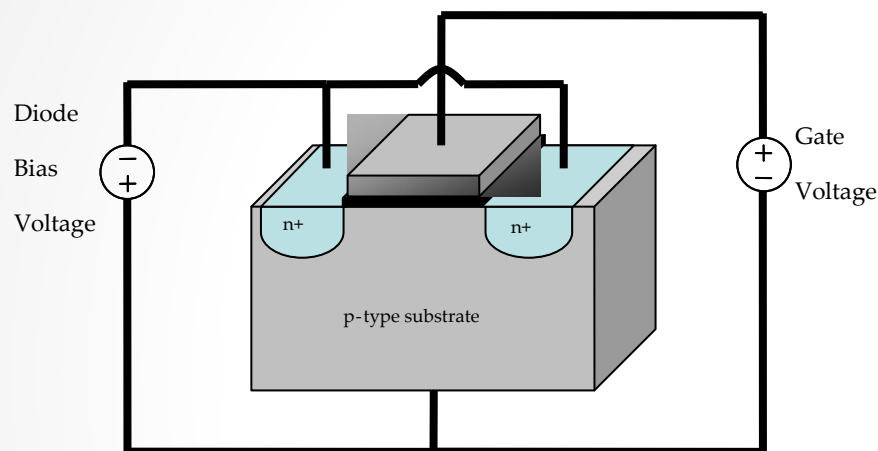
Nature Materials 9, 725 (2010)



What about Read-out ?

- Long relaxation times
- Initialization
- Manipulation
- Readout ??
 - Read-out is usually destructive
 - Magnetic resonance is not very sensitive 10^9 spins

Detection



Electrical detection of the spin resonance of a single electron in a silicon field-effect transistor

M. Xiao¹, I. Martin², E. Yablonovitch³ & H. W. Jiang¹

¹Department of Physics and Astronomy, University of California, Los Angeles, California 90095, USA

²Theoretical Division, Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA

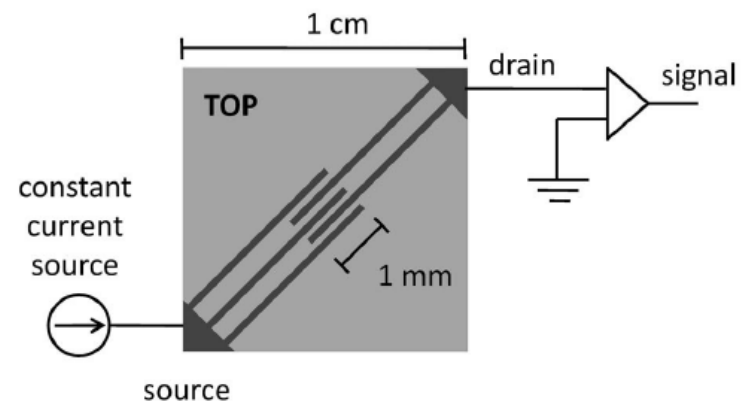
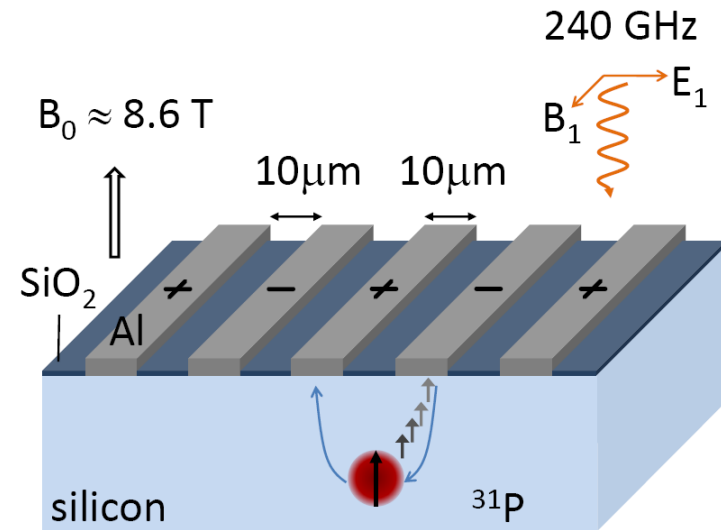
³Department of Electrical Engineering, University of California, Los Angeles, California 90095-1594, USA

MOSFET : spin-dependent recombination

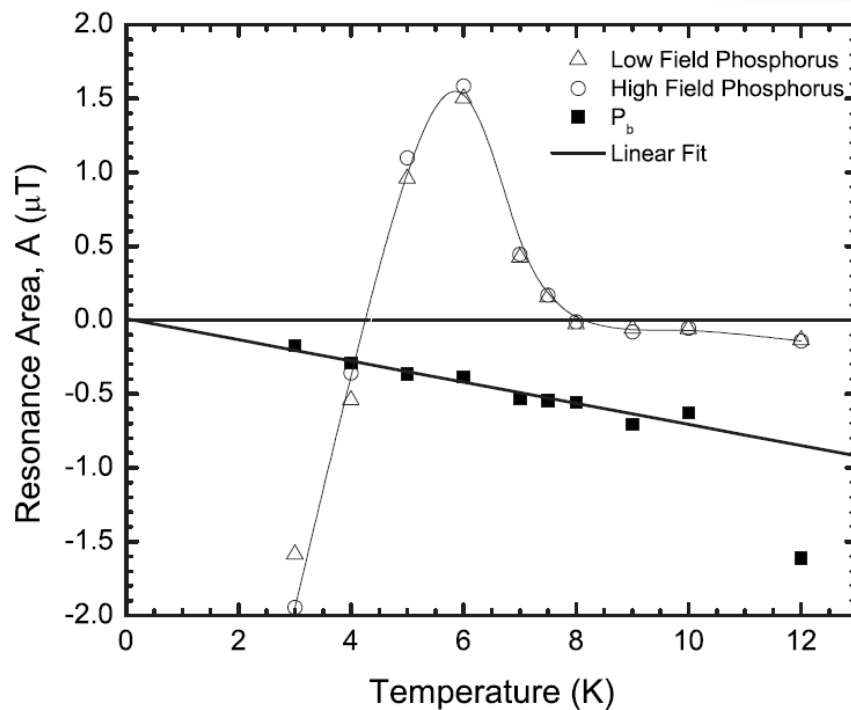
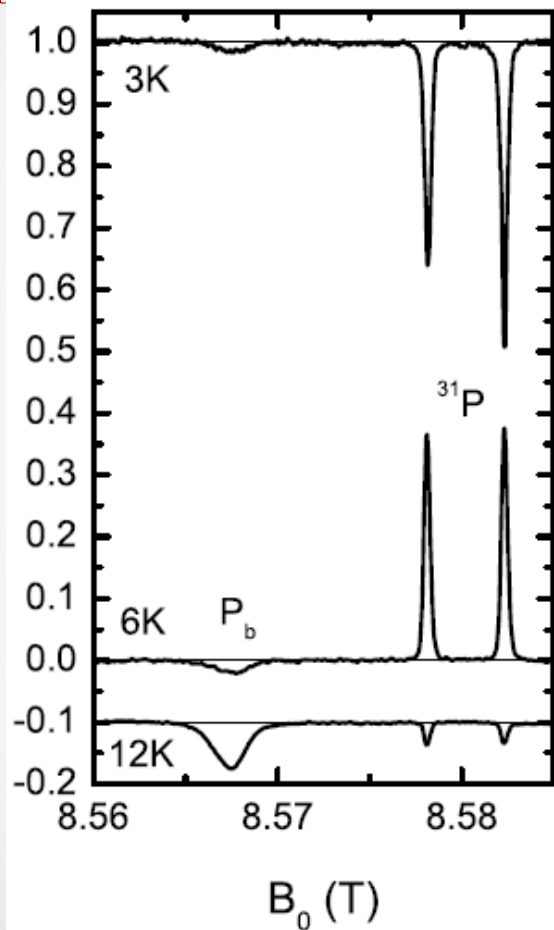
High frequency advantage : Boltzmann populations

Electrical detection of Magnetic Resonance (Read out)

- Electrical detection can be very sensitive
- At low frequencies a P_b (surface) center is involved
 - In order to measure a current we use light excitation to create carriers
 - T_1 is shortened, T_2 more or less unchanged



EDMR spectra

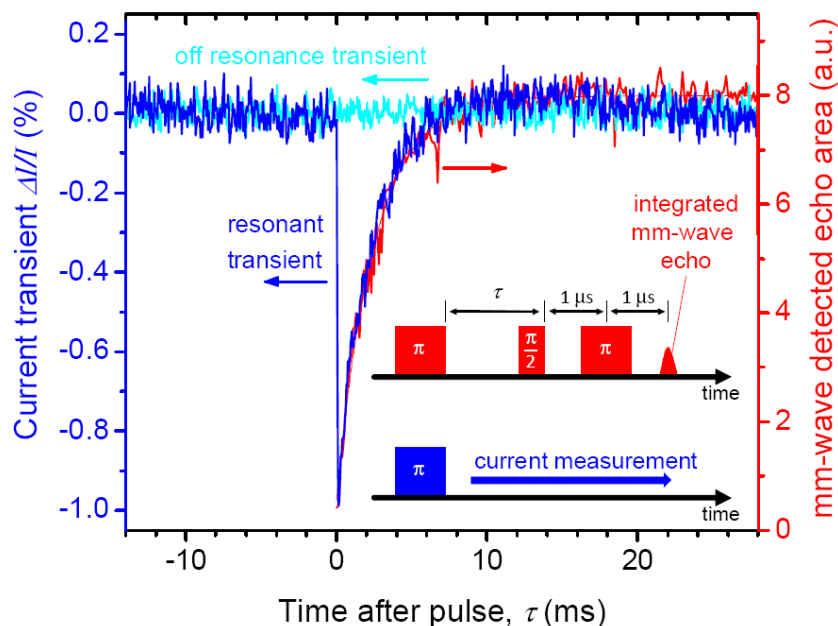


Up to 10% current changes on
phosphorus resonance

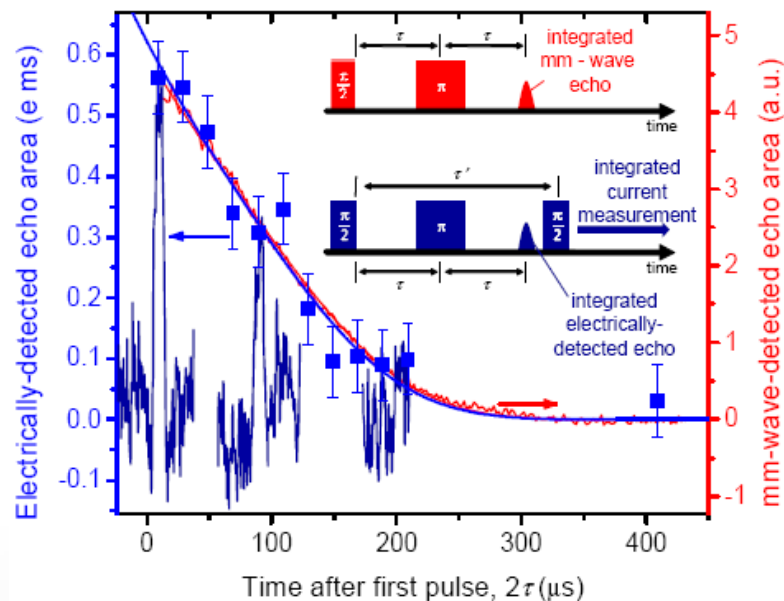
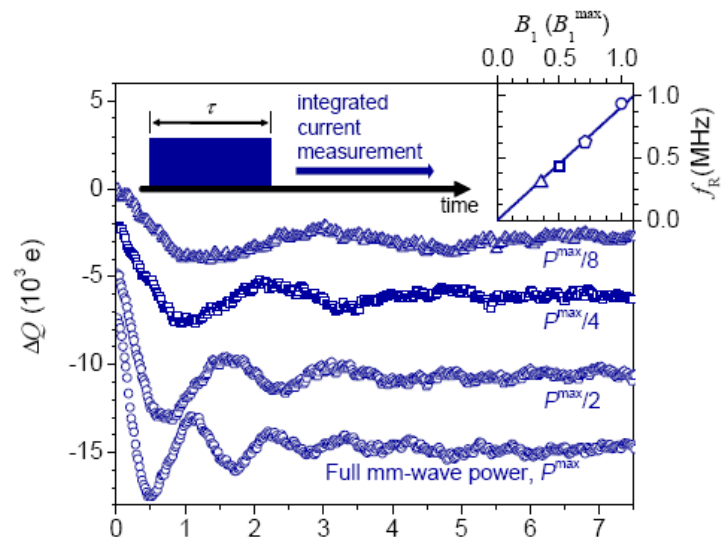
McCamey et al., PRB 2008

Coherence ?

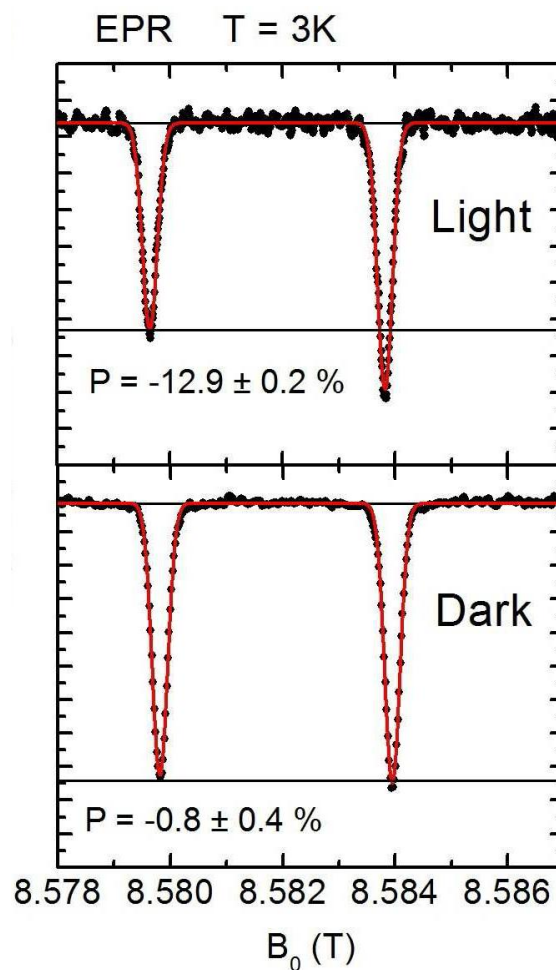
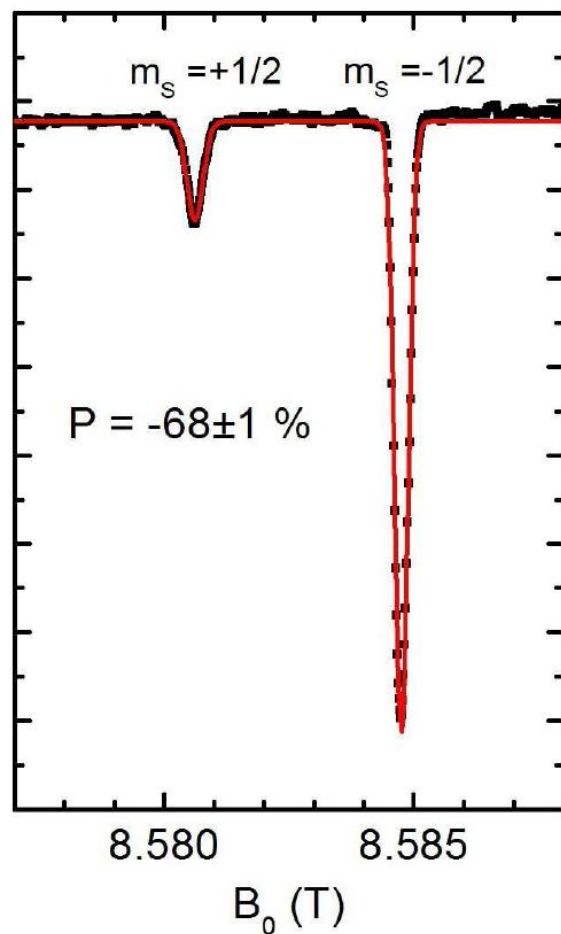
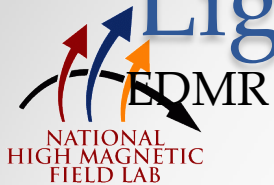
- At low frequency the electrically detected spin coherence lasts for $\sim 2 \mu\text{s}$ due to fast recombination of electron-hole pairs.
- At high frequency the main mechanism is dominated by spin traps $\rightarrow$ longer coherence times ?



Morley et al, PRL 101 (2008)

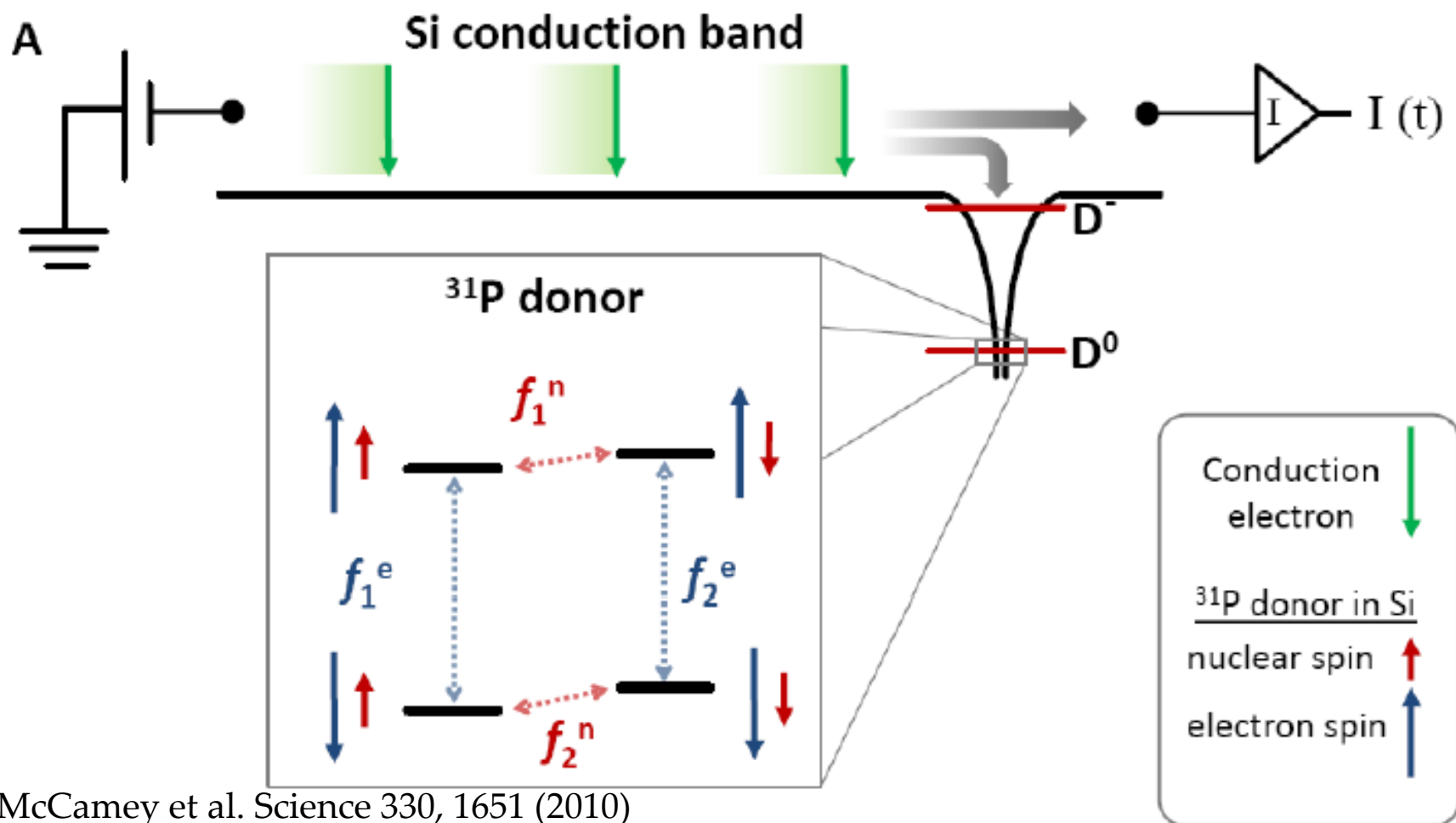


Light Induced Nuclear Polarization

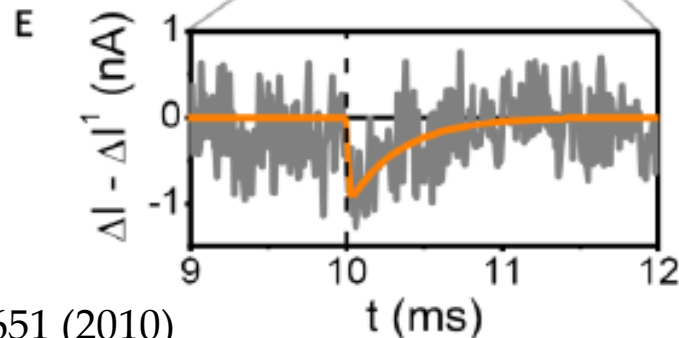
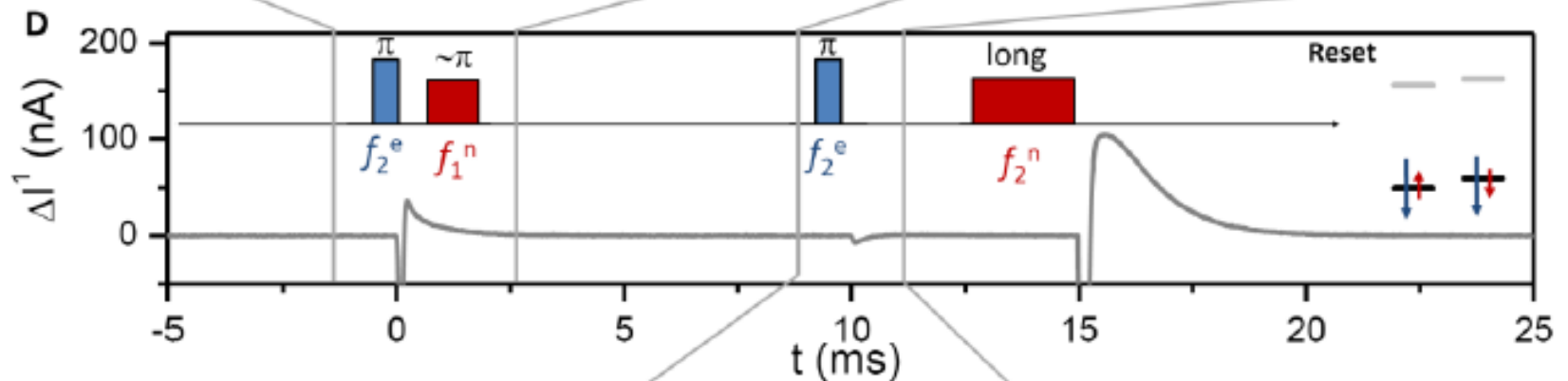
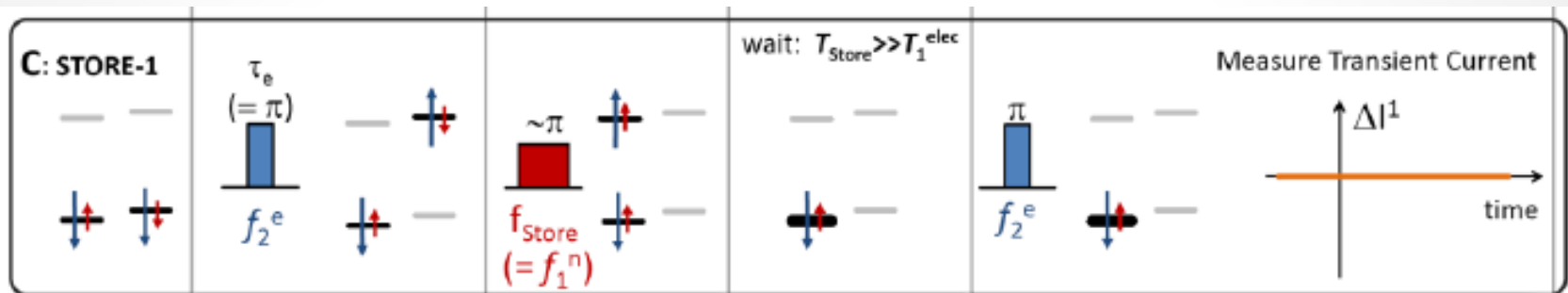


Storage of information

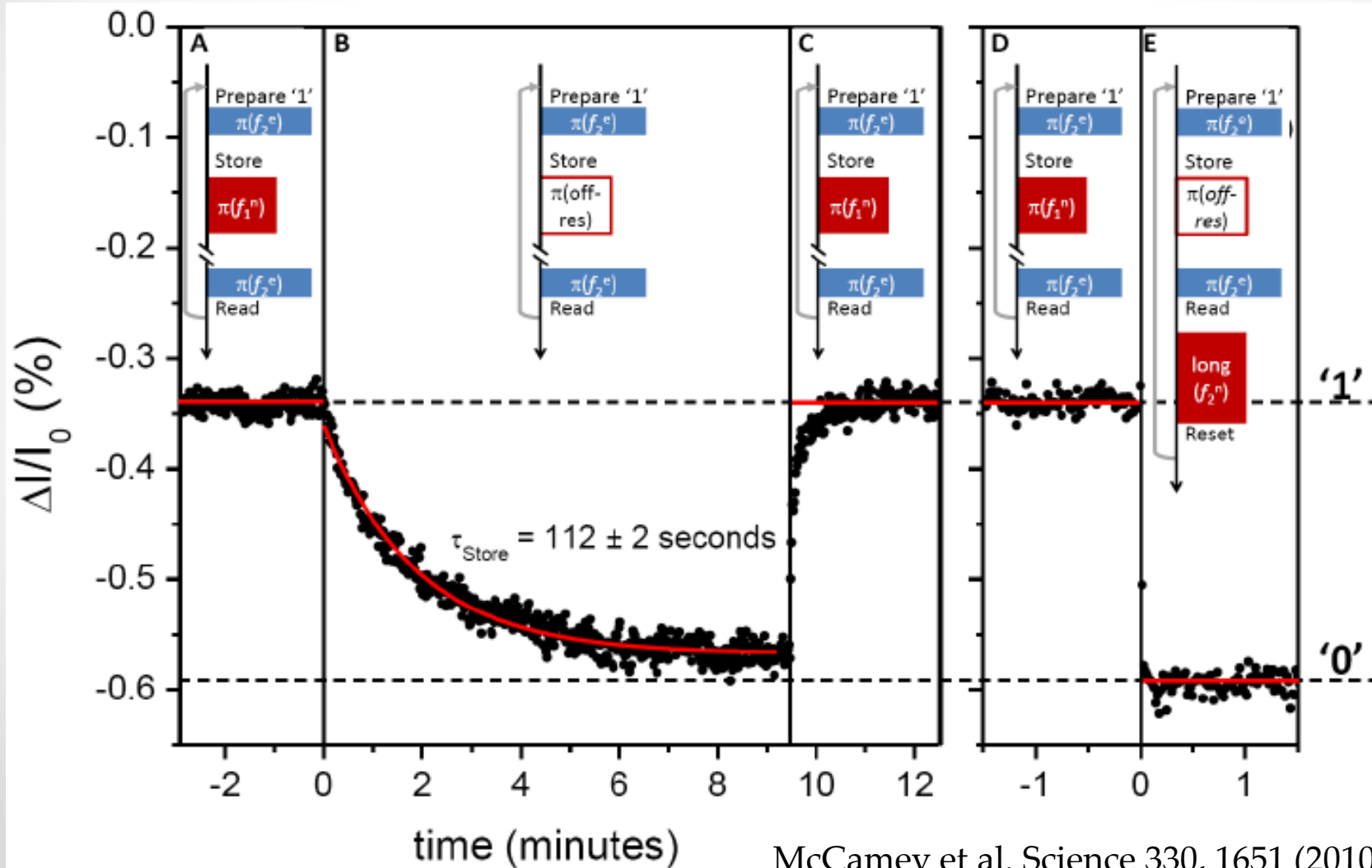
- Use the nuclear spin to store electron spin



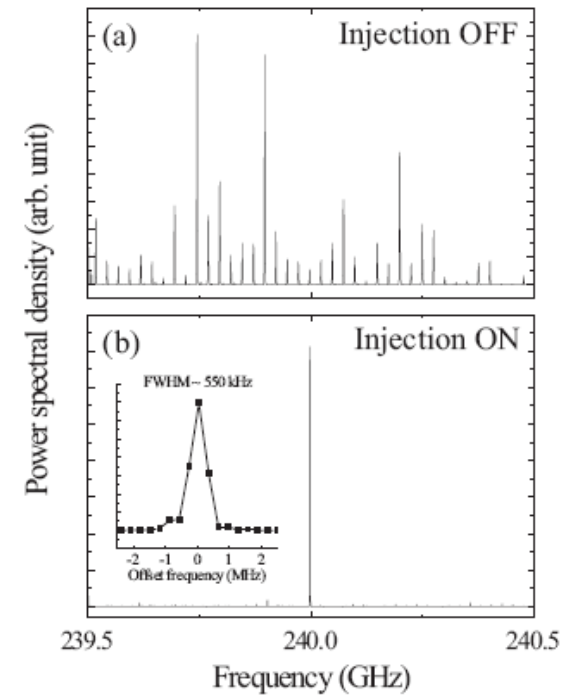
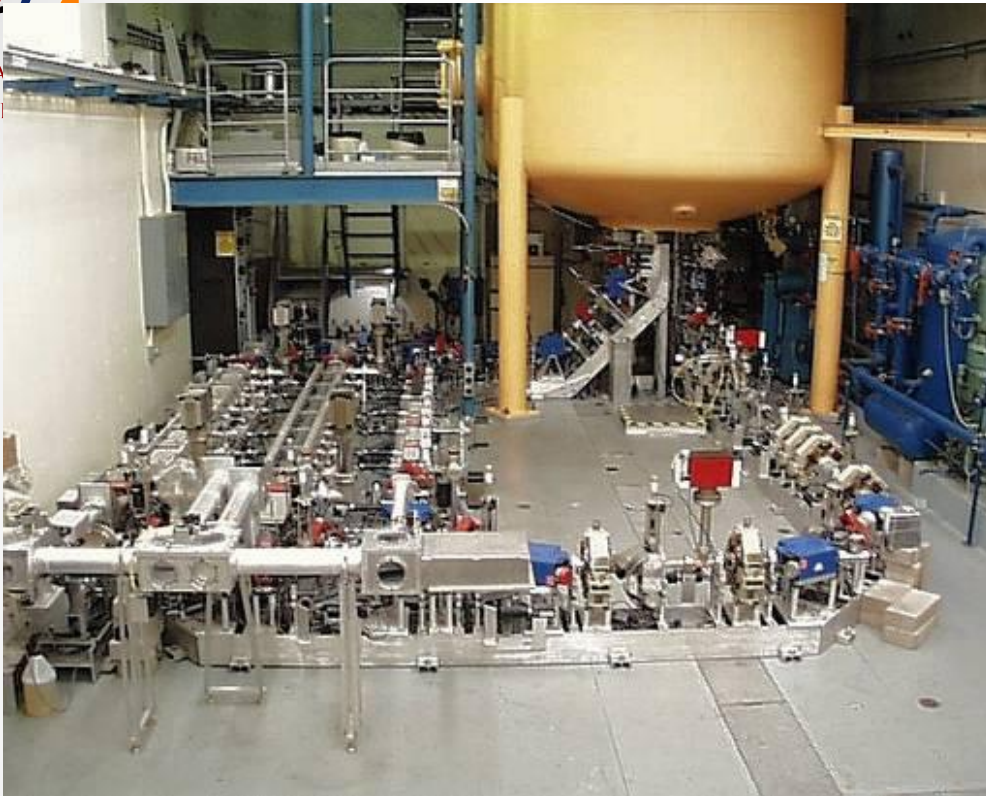
Electrical storage and readout of spin information



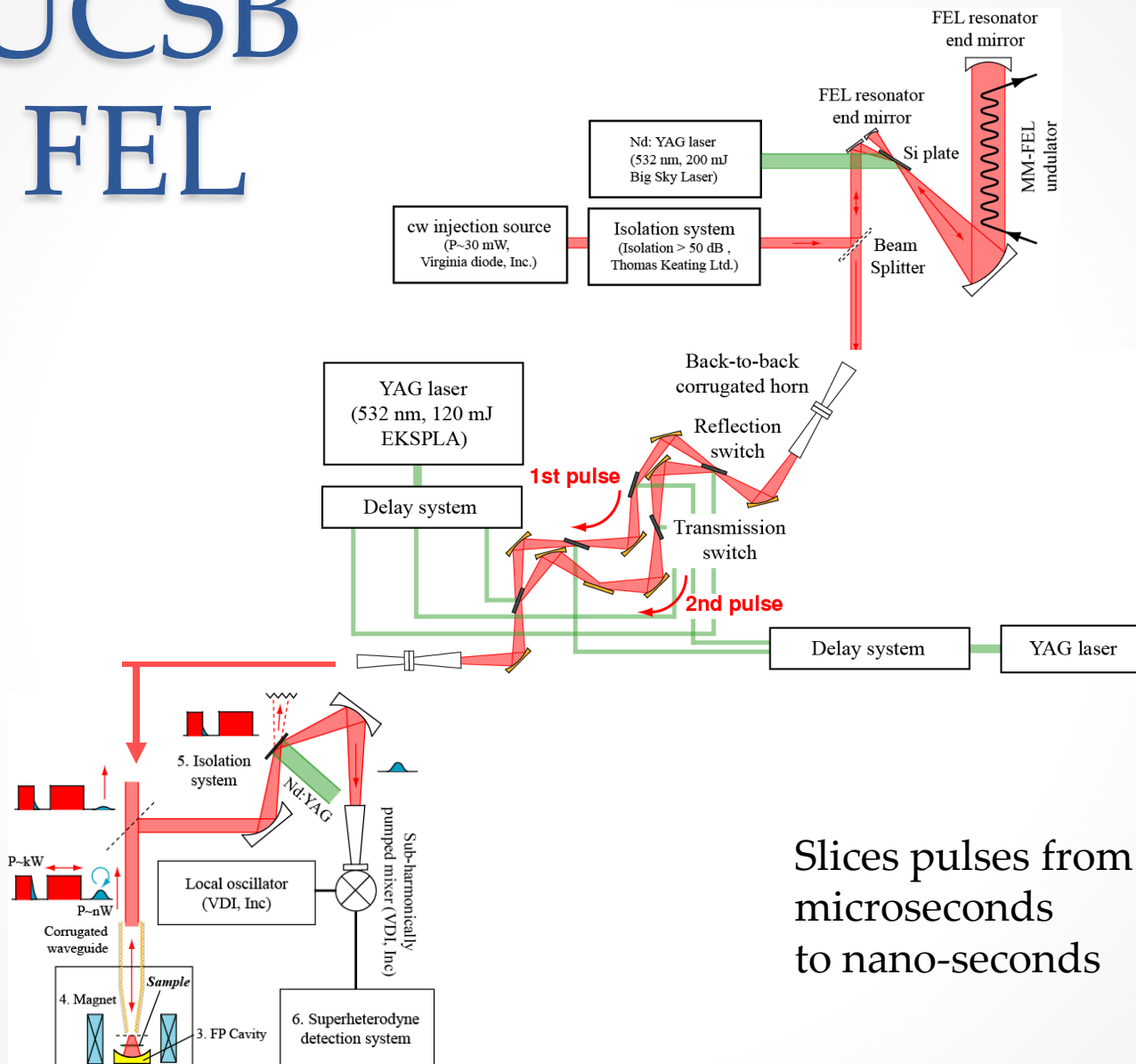
A minute or so...



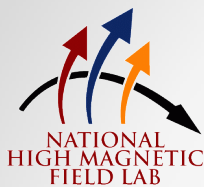
How strong is the spin dependent current



UCSB FEL



Slices pulses from
microseconds
to nano-seconds



Acknowledgements

Gavin Morley
Dane McCamey
Christoph Boehme

UCL (London)
U. Utah (U Sidney)
U. Utah/Phys

Susumu Takahashi
Mark Sherwin
Louis-Claude Brunel
Devin Edwards

UCSB (USC)
UCSB/ITST
UCSB/ITST
UCSB Chem

